

# Recent developments in interval dynamics

Weixiao Shen and Sebastian van Strien

**Abstract.** Dynamics in dimension-one has been an extremely active research area over the last decades. In this note we will describe some of the new developments of the recent years.

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## 1. Density of hyperbolicity

Interval maps  $f: [0, 1] \rightarrow [0, 1]$  can have a surprisingly rich and complicated dynamics. In this paper we will describe results which show that in spite of this one can describe the metric orbit structure of ‘most’ maps extremely well.

The dynamics of *hyperbolic* maps can be described most easily: for these maps, Lebesgue almost every point in the interval is attracted to some hyperbolic periodic orbit (with multiplier between  $-1$  and  $1$ ). By a result by Mañé [65] (for a simpler proof see [98]) it is equivalent to say that a map is hyperbolic if (i) each critical point of  $f$  is in the basin of a periodic attractor and (ii) each periodic orbit is hyperbolic. Since the period of periodic attractors is bounded, see [66], it follows that hyperbolic maps have at most finitely many periodic attractors.

As mentioned, hyperbolic maps are very well-understood. The following theorem (which was obtained by the authors, jointly with Kozlovski, see [50]) shows that ‘most’ maps are hyperbolic.

**Theorem 1.1** (Density of hyperbolicity for real polynomials). *Any real polynomial can be approximated by hyperbolic real polynomials of the same degree.*

The above theorem allows us to prove the analogue of the Fatou conjecture in the smooth case, see [51], thus solving the 2nd part of Smale’s eleventh problem for the 21st century [91]:

**Theorem 1.2** (Density of hyperbolicity for smooth one-dimensional maps). *Hyperbolic maps are dense in the space of  $C^k$  maps of the compact interval or the circle,  $k = 1, 2, \dots, \infty, \omega$ .*

For quadratic maps  $f_a = ax(1-x)$ , the above theorems assert that the periodic windows (corresponding to hyperbolic maps with attracting periodic orbits) are dense in the bifurcation diagram. The quadratic case turns out to be special, because in this case certain return maps become almost linear. This special behaviour does not even hold for maps of the form  $x \mapsto x^4 + c$ .

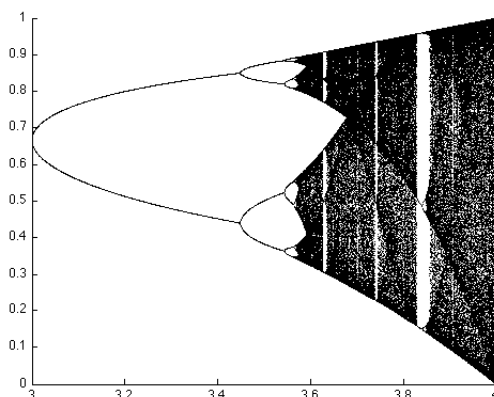


Figure 1.1. The Feigenbaum diagram

35 The problem of density of hyperbolicity in dimension-one has been considered since the  
 36 1920's. Indeed:

- 37 • Fatou stated the analogue of this problem in the context of rational maps on the Riemann  
 38 sphere as a conjecture in the 1920's, see [33, page 73] and also [67, Section 4.1].
- 39 • Smale gave this problem 'naively' as a thesis problem in the 1960's, see [90].
- 40 • In 1971, Jakobson proved that the set of hyperbolic maps is dense in the  $C^1$  topology, see  
 41 [43].
- 42 • In the mid 1990's, the conjecture was solved in the quadratic case  $x \mapsto ax(1-x)$  in  
 43 a major breakthrough by Lyubich [61] and independently also by Graczyk and Świątek,  
 44 [36] and [37].
- 45 • In 2000, Blokh and Misiurewicz [15] considered the problem of density of hyperbolicity  
 46 in the  $C^2$  topology, and were able to obtain a partial result.
- 47 • A few years later, Shen [87] proved  $C^2$  density of hyperbolic maps.

48 Note that every hyperbolic map satisfying a mild transversality condition, namely that  
 49 no critical point is eventually mapped onto another critical point, is *structurally stable*. So  
 50 density of hyperbolicity implies that structural stable maps are dense.

51 **1.1. Density of hyperbolicity within a large space of real transcendental map.** Density  
 52 of hyperbolicity also holds within classes of much more general maps, for example within  
 53 the famous Arnol'd family and within the space of trigonometric polynomials. Indeed it was  
 54 shown by the second author in a joint paper with Rempe, see [79], that

55 **Theorem 1.3.** *Density of hyperbolicity holds within the following spaces:*

- 56 1. *real transcendental entire functions, bounded on the real line, whose singular set is*  
 57 *finite and real;*
- 58 2. *transcendental functions  $f: \mathbb{C} \setminus \{0\} \rightarrow \mathbb{C} \setminus \{0\}$  that preserve the circle and whose*  
 59 *singular set (apart from  $0, \infty$ ) is contained in the circle.*

In [79] also a number of other open problems are solved, including a number of conjectures of behaviour de Melo, Salomão and Vargas [29]. In this paper density of (real) hyperbolicity is also established replacing in assumption (1) the boundedness condition by a sector condition.

## 1.2. Hyperbolicity is dense within generic one-parameter families of one-dimensional maps.

**Theorem 1.4** (Hyperbolicity is dense within generic families). *For any generic family  $\{g_t\}_{t \in [0,1]}$  of smooth interval maps (generic, in the sense of Baire), the following properties hold:*

- *the number of critical points of each of the maps  $g_t$  is bounded;*
- *the set of parameters  $t$  for which all critical points of  $g_t$  are in basins of periodic attractors, is open dense.*

The proof of this result follows easily from the theorems in the previous subsection, see [99]. On the other hand, as is shown in the same paper, it is easy to construct a real analytic one-parameter family  $f_t$ ,  $t \in [0, 1]$  of polynomials so that none of the polynomials in this family are hyperbolic:

**Theorem 1.5** (A family of cubic maps with robust chaos). *There exists a real analytic one-parameter family  $\{f_t\}$  of interval maps (consisting of cubic polynomials) so that  $f_t$  has no periodic attractor for any  $t \in [0, 1]$ , and so that not all maps within this family are topologically conjugate.*

**1.3. Density of hyperbolicity for more general maps.** Density of hyperbolicity is false in dimension  $\geq 2$ . For a list of related interesting questions concerning the higher dimensional case, see [77].

The situation for rational maps on the Riemann sphere may well be more hopeful. In that context one has the following well-known conjecture, going back to Fatou:

**Conjecture 1.6** (Density of hyperbolicity for rational maps). *Hyperbolic maps are dense within this space of rational maps of degree  $d$  on the Riemann sphere.*

In [64] it was shown that this conjecture follows from

**Conjecture 1.7.** *If a rational map carries a measurable invariant line field on its Julia set, then it is a Lattès map.*

More about this conjecture and related results can be found in [67]. In [50, 86] and finally [52] it was shown that real polynomials (acting on  $\mathbb{C}$ ) do not carry such invariant line fields. Moreover, real polynomials have Julia sets which are locally connected, see [26, 50, 52, 55]. In [80] it was shown that, under some mild assumptions, real transcendental maps also do not carry invariant line fields.

Interestingly, any rational map on the Riemann sphere such that the multiplier of each periodic orbit is real, either has a Julia set which is contained in a circle (or line) or is a Lattès map, see [32].

**1.4. Strategy of the proof: local versus global perturbations.** Density of hyperbolicity means that given a map  $f$  one can find a map  $g$  so that  $g$  is hyperbolic and so that  $g -$

100  $f$  is ‘small’ in the  $C^k$  topology. It is tempting to consider the setting where  $g$  is a local  
 101 perturbation of  $f$ . The purpose would then be to find a small ‘bump’ function  $h$  so that  
 102  $g = f + h$  becomes hyperbolic. The difficulty with this approach is that orbits will pass  
 103 many times through the support of the bump function. Pugh’s approach in his proof of the  
 104  $C^1$  closing lemma, is to find a suitable neighbourhood  $U$  of  $x$  so that the first return of  $x$  to  
 105  $U$  is not too close to the boundary of  $x$ . In this way he is able to construct a function  $h$  whose  
 106 support is in  $U$ , which creates a new fixed point of the first return of  $g = f + h$  to  $U$ , in such  
 107 a manner that  $h$  is  $C^1$  close to zero. A related approach was used successfully in [43] to  
 108 prove density of hyperbolicity in the  $C^1$  topology, and in [15] for the  $C^2$  topology, but with  
 109 added assumptions on the dynamics of  $f$ . In [87], this approach was used in the case when  
 110 one has a ‘lot of Koebe space’ while in the ‘essentially bounded geometry’ the proof relied  
 111 on rigidity (in the sense described below). This rigidity approach also is the key ingredient  
 112 in the proof of Theorem 1.1. As there is a great deal of evidence that local perturbations  
 113 cannot be used to prove density of hyperbolicity in general, we discuss rigidity extensively  
 114 in the next section.

115 **1.5. Strategy of the proof: quasi-symmetric rigidity.** Consider the following situation.  
 116 Take a family of real quadratic maps  $f_c(z) = z^2 + c$ . To prove density of hyperbolicity  
 117 we need to prove that there exists no interval of parameters  $[c', c'']$  so that each map  $f_c$   
 118 with  $c \in [c', c'']$  is non-hyperbolic. Sullivan showed that this follows from quasi-symmetric  
 119 rigidity of any non-hyperbolic map  $f_c$ . Here  $f_c$  is called *quasi-symmetrically rigid* if the  
 120 following property holds:

121 If  $f_{\tilde{c}}, f_{\hat{c}}$  are topologically conjugate to  $f_c$ , then  $f_{\tilde{c}}, f_{\hat{c}}$  are quasi-symmetrically  
 122 conjugate.

Here, as usual, a homeomorphism  $h: [0, 1] \rightarrow [0, 1]$  is called *quasi-symmetric* (often abbrevi-  
 ated as *qs*) if there exists  $K < \infty$  so that

$$\frac{1}{K} \leq \frac{h(x+t) - h(x)}{h(x) - h(x-t)} \leq K$$

123 for all  $x - t, x, x + t \in [0, 1]$ . By results about quasi-conformal maps (specifically the  
 124 Measurable Riemann Mapping Theorem) it follows that the set of parameters  $\tilde{c}$  so that  $f_{\tilde{c}}$  is  
 125 topologically conjugate to  $f_c$  is either a single point or an open interval  $I(f_c)$ . Since  $I(f_c)$  is  
 126 also a closed set (this follows from some basic kneading theory), the fact that  $I(f_c)$  and its  
 127 complement are both non-empty gives a contradiction unless  $I(f_c)$  is a single point.

128 This argument does not go through directly for real polynomial maps with more than  
 129 one critical point, but using related arguments, one still obtains that quasi-symmetric rigidity  
 130 implies density of hyperbolicity, see [50, Section 2]. In the case of real analytic maps the  
 131 argument to prove density of hyperbolicity is more subtle, see [51].

## 132 2. Quasi-symmetric rigidity

133 As remarked in the previous section, all current proofs of density of hyperbolicity rely on  
 134 quasi-symmetric rigidity. The most general form can be found in [25], and states:

135 **Theorem 2.1** (Quasi-symmetric rigidity). Assume that  $f, g: [0, 1] \rightarrow [0, 1]$  are real analytic  
 136 and topologically conjugate. Alternatively, assume that  $f, g: S^1 \rightarrow S^1$  are topologically  
 137 conjugate and that  $f$  and  $g$  each have at least one critical point or at least one periodic  
 138 point. Moreover, assume that the topological conjugacy is a bijection between

- 139 (1) the set of critical points and the order of corresponding critical points is the same;  
 140 (2) the set of parabolic periodic points.

141 Then the conjugacy between  $f$  and  $g$  is quasi-symmetric.

142 The proof of this theorem builds on the machinery developed in [50]. This paper was  
 143 written jointly by the authors and Kozlovski; it developed many of the key ingredients re-  
 144 quired to prove density of hyperbolicity, see [51]. Theorem 2.1 is an extension of these  
 145 results, and was obtained jointly by Clark and the 2nd author, and uses all of the technology  
 146 from [51], but also extends ideas from [56].

147 Indeed, when  $f, g$  are real analytic, then we will use the fact that these maps have holo-  
 148 morphic extensions to small neighbourhoods of  $[0, 1]$ . Nevertheless, in [25] we prove the  
 149 analogous result when  $f$  and  $g$  are merely  $C^3$  maps, under some weak additional assump-  
 150 tions; in this case we will use that  $f, g$  have asymptotically holomorphic extensions near  
 151  $[0, 1]$ , but will need to deal with the fact that high iterates of  $f$  and  $g$  are not necessarily close  
 152 to holomorphic.

153 It is not hard to see that if conditions (1) or (2) in the previous theorem are not satisfied,  
 154 then the maps are not even necessarily Hölder conjugate.

155 Special cases of this theorem we known before: Lyubich [61] and Graczyk & Świątek  
 156 [37] proved this result for real quadratic maps. As we will see their method of proof in  
 157 the quadratic case does not work if the degree of the map is  $> 2$ . For the case of real  
 158 polynomials with only real critical points (of even order), this theorem was proved in [50].  
 159 For maps which are real analytic, it was shown in [87, Theorem 2, page 345] that there  
 160 exists a qs-conjugacy *restricted to*  $\omega(c)$  under the additional assumptions that the maps have  
 161 no neutral periodic points, only non-degenerate critical points and have ‘*essentially bounded*  
 162 *geometry*’. For covering maps of the circle (of degree  $\geq 2$ ) with one-critical point a global  
 163 qs-conjugacy was constructed under the additional assumption that  $\omega(c)$  is non-minimal and  
 164 have no neutral periodic points, see [56]. When  $\omega(c)$  is minimal, a qs-conjugacy restricted  
 165 to  $\omega(c)$  was constructed in [56].

166 For circle maps without periodic points, it is known that any two analytic critical circle  
 167 homeomorphisms with one critical point, with the same irrational rotation number and the  
 168 same order of the critical points are  $C^1$ -smoothly conjugate, see [47] (their work builds on  
 169 earlier work of de Faria, de Melo and Yampolsky on renormalisation and in a recent paper  
 170 was generalised to the smooth case, [39]). In ongoing work, Clark and the 2nd author are  
 171 aiming to show that the methods in 2.1 can be extended to the case of circle homeomorphisms  
 172 with several critical points. Note that the presence of critical points is necessary for circle  
 173 homeomorphisms, because for circle diffeomorphisms the analogous statement is false. In-  
 174 deed, otherwise one can construct maps for which some sequence of iterates has *almost a*  
 175 *saddle-node fixed point*, resulting in larger and larger passing times near these points. This  
 176 phenomenon is also referred to as *a sequence of saddle-cascades*. It was used by Arnol’d  
 177 and Herman to construct examples of diffeomorphisms of the circle which are conjugate to  
 178 irrational rotations, but where the conjugacy is neither absolutely continuous, nor qs and for  
 179 which the map has no  $\sigma$ -finite measures, see [40] and also Section I.5 in [30]. In the diffeo-

180 morphic case, to get  $qs$  or  $C^1$  one needs assumptions on the rotation number (to avoid these  
181 sequences of longer and longer saddle-cascades).

182 In general, one cannot expect  $C^1$ , because having a  $C^1$  conjugacy implies that corre-  
183 sponding periodic orbits have the same multiplier.

184 We should also remark that there are also analogues of these theorems for polynomials  
185 in  $\mathbb{C}$ , but then one must assume that  $f$  is only finitely renormalizable, see for example [52],  
186 but also see [24].

187 **2.1. Applications of quasi-symmetric rigidity.** Quasi-symmetric rigidity is a crucial step  
188 towards proving the following types of results:

- 189 (1) hyperbolicity is dense, see subsection 1.5.
- 190 (2) within certain families of maps, conjugacy classes are connected, see Theorems A and  
191 2.2 in [21].
- 192 (3) monotonicity of entropy; for families such as  $[0, 1] \ni x \mapsto a \sin(\pi x)$ , see Section 3.

193 **2.2. Complex box mappings.** It turns out to be rather convenient to show quasi-symmetric  
194 rigidity by using extensions to the complex plane. This approach is rather natural, as a quasi-  
195 symmetric homeomorphism on the real line is always the restriction of a quasi-conformal  
196 homeomorphism on the complex plane. More precisely, the idea is to construct an extension  
197 of the first return map to some interval, to the complex plane as a ‘complex box mapping’, see  
198 Figure 2.1 in the multimodal case. Roughly speaking, this is a map  $F: U \rightarrow V$  so that each  
199 component of  $U$  is mapped as a branched covering onto a component of  $V$ , and components  
200 of  $U$  are either compactly contained in a component of  $V$  or they are equal to such a compo-  
201 nent. Components of  $F^{-n}(V)$  are called *puzzle pieces*. We also require (roughly speaking)  
202 that  $F$  is unbranched near the boundary of  $U$  (slightly more precisely, that there exists an  
203 annulus neighbourhood  $A$  of  $\partial V$  so that  $F: F^{-1}(A) \rightarrow A$  is an unbranched covering and so  
204 that  $\text{mod}(A)$  is universally bounded from below). If one has such numerical bounds, then  $F$   
205 is said to have *complex bounds*. The existence of these complex bounds was first proved by  
206 Sullivan for certain unimodal maps. The general unimodal case was dealt with in [55] and  
207 somewhat later in [35] and [60]. Later this was extended to the multimodal case for certain  
208 maps in [92] and more generally in [87]. The most general result appears in a joint paper  
209 of the 2nd author with Clark and Trejo [26]. In that paper complex bounds are associated to  
210 any real analytic interval map. In fact, even in the  $C^3$  case complex bounds are constructed  
211 in that paper, but in the smooth case the map  $F$  is only asymptotically holomorphic.

212 We should note that in the non-renormalisable real-analytic case one obtains complex  
213 bounds at arbitrary deep levels, as soon as one has a complex box mapping. That this is the  
214 case follows from the construction of the enhanced nest (discussed in the next subsection)  
215 and an interesting lemma due to Kahn and Lyubich, see [45]. This tool is about pulling back  
216 a thin annulus, and shows that the modulus of the pullback of this annulus is much better  
217 than one might expect. In the real case, one can simplify and strengthen the statement and  
218 proof of Kahn and Lyubich’s result as follows, see [52, Lemma 9.1]:

**Lemma 2.2** (Small Distortion of Thin Annuli). *For every  $K \in (0, 1)$  there exists  $\kappa > 0$   
such that if  $A \subset U$ ,  $B \subset V$  are simply connected domains symmetric with respect to the real  
line,  $F: U \rightarrow V$  is a real holomorphic branched covering map of degree  $D$  with all critical  
points real which can be decomposed as a composition of maps  $F = f_1 \circ \dots \circ f_n$  with all  
maps  $f_i$  real and either real univalent or real branched covering maps with just one critical*

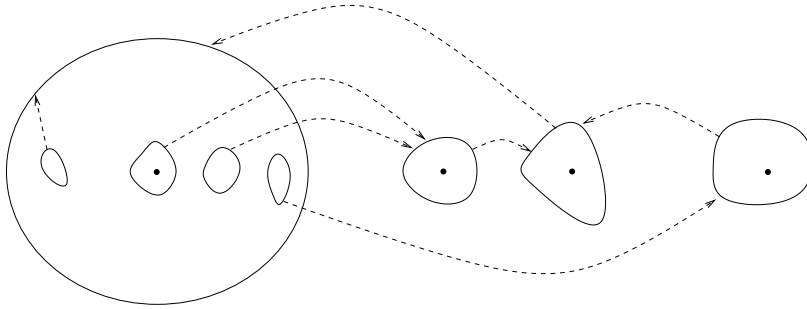


Figure 2.1. A box mapping.

point, the domain  $A$  is a connected component of  $f^{-1}(B)$  symmetric with respect to the real line and the degree of  $F|_A$  is  $d$ , then

$$\text{mod}(U - A) \geq \frac{K^D}{2d} \min\{\kappa, \text{mod}(V - B)\}.$$

219 **2.3. How to prove quasi-symmetric rigidity?** Consider the complex box mappings asso-  
 220 ciated to two conjugate maps. To show that the conjugacy is quasi-symmetric one proceeds  
 221 as follows:

- 222 (1) Define a sequence of puzzle pieces  $U_{n_i}$  called the **enhanced nest**, so that there exists  
 223  $k_i$  for which  $F^{k(i)}: U_{n_{i+1}} \rightarrow U_{n(i)}$  is a branched covering map with degree bounded  
 224 by some universal number  $N$ . This enhanced nest is chosen so that it transfers geomet-  
 225 ric information rather efficiently from small scale to large scale, but so that the degree  
 226 of  $F^{k(i)}: U_{n_{i+1}} \rightarrow U_{n(i)}$  remains universally bounded. This enhanced nest was one  
 227 of the main new ingredients in [50]. It turns out that the post-critical sets do not come  
 228 close to the boundary of the puzzle pieces in the enhanced nest, which implies that the  
 229 puzzle pieces have uniformly bounded shape. Another important property of the en-  
 230 hanced nest is that decaying geometry and bounded geometry alternate quite regularly  
 231 in the nest, which was used in [58] to study the Hausdorff dimension of Cantor attrac-  
 232 tors. The enhanced nest construction is also used for example in [26, 52, 75, 78, 94].
- 233 (2) In fact, if the interval maps extend to a holomorphic map on a neighbourhood of the  
 234 real line, then one can partially define a quasi-conformal conjugacy near critical points,  
 235 and then spread the definition to the whole complex plane fairly easily. This method  
 236 was called the *spreading principle* in [50].
- 237 (3) Because of the spreading principle mentioned above, it then suffices to construct a  
 238 partial-conjugacy on a puzzle piece in the enhanced nest which is ‘natural on the  
 239 boundary’. Given the above, this can easily be done using the *QC-criterion* from  
 240 the appendix of [50]. and bounded shape of the puzzle pieces (bounded shape is very  
 241 easy to derive from complex bounds, see [52, Section 10]). One can also proceed as in  
 242 [5]. Our QC criterion was a variation of Heinonen-Koskela’s theorem [42]. This theo-  
 243 rem and its variations were used to prove rigidity result previously in [34, 41, 56, 83?  
 244 ], where in the last work, the author explicitly stated that a bounded shape property of  
 245 puzzle pieces implies rigidity for non-renormalizable unicritical maps.

246 It is of course conceivable that one can prove quasi-symmetric rigidity using entirely real  
247 methods. This hinges on questions of the following type:

248 **Question 2.3.** *Consider the space  $\mathcal{A}$  of maps of the form  $z \mapsto |z|^d + c$  where  $d > 1$  is not  
249 necessarily an integer and where  $c$  is real. Does one have quasi-symmetric rigidity for maps  
250 within the space  $\mathcal{A}$ ? Are two topologically conjugate maps in  $\mathcal{A}$  without periodic attractors  
251 (or both critically finite) necessarily the same?*

252 One of the difficulties with such a real approach is that it is not so easy to know how  
253 to use the information that the exponent  $d$  is fixed within the family  $\mathcal{A}$ : the exponent is not  
254 ‘visible’ in the real line. On the other hand, if  $d$  is an even integer, and  $z \mapsto z^d + c$ , then of  
255 course the local degree of the map at 0 is different for different values of  $d$ . Without fixing  
256 the degree  $d$  the answer to the question above is definitely negative. An affirmative answer  
257 to the above question would imply density of hyperbolicity and monotonicity of entropy in  
258 this family.

### 259 3. Monotonicity of entropy

260 In the late 70’s, the following question attracted a lot of interest: does the topological entropy  
261 of the interval map  $x \mapsto ax(1-x)$  depend monotonically on  $a \in [0, 4]$ ? In the mid 80’s this  
262 question was solved in the affirmative:

263 **Theorem 3.1.** *The topological entropy of the interval map  $x \mapsto ax(1-x)$  depends mono-  
264 tonically on  $a \in [0, 4]$ .*

265 In the 80’s several proofs of this appeared. One of these uses Thurston’s rigidity theorem,  
266 see [70]. Another proof relies on Douady-Hubbard’s univalent parametrisation of hyperbolic  
267 components, see [31], and a third proof is due to Sullivan; for a description of these proofs  
268 see [30]. All these proofs consider the map  $x \mapsto ax(1-x)$  as a polynomial acting on the  
269 complex plane. A rather different method was used by Tsujii, [97]. He showed that periodic  
270 orbits bifurcate in the ‘right’ direction using a calculation on how the multiplier depends  
271 on the parameter. Unfortunately, Tsujii’s proof also does not work for maps of the form  
272  $z \mapsto |z|^a + c$  with  $a$  not an integer.

273 In the early 90’s, Milnor (see [69]) posed the more general

274 **Conjecture 3.2** (Monotonicity Conjecture). *The set of parameters within a family of real  
275 polynomial interval maps, for which the topological entropy is constant, is connected.*

276 Milnor and Tresser proved this conjecture for cubic polynomials, see [71] (see also  
277 [28]). Their ingredients are planar topology (in the cubic case the parameter space is two-  
278 dimensional) and density of hyperbolicity for real quadratic maps.

279 A few years ago, Bruin and the 2nd author were able to give a proof of the general case  
280 of this conjecture. More precisely, given  $d \geq 1$  and  $\epsilon \in \{-1, 1\}$ , consider the space  $P_\epsilon^d$  of  
281 real polynomials  $f: [0, 1] \rightarrow [0, 1]$  of fixed degree  $d$  with  $f(\{0, 1\}) \subset \{0, 1\}$ , with all critical  
282 points in  $(0, 1)$  and with the first lap orientation preserving if  $\epsilon = 1$  and orientation reversing  
283 if  $\epsilon = -1$ . We call  $\epsilon$  the *shape* of  $f$ . In [21] we proved the general case:

**Theorem 3.3** (Monotonicity of Entropy). *For each integer  $d \geq 1$ , each  $\epsilon \in \{-1, 1\}$  and  
each  $c \geq 0$ ,*

$$\{f \in P_\epsilon^d; h_{top}(f) = c\}$$



284 *is connected.*

285 The proof in [21] also shows that the set of maps in  $P_\epsilon^d$  with the same kneading sequence  
 286 is connected and gives a precise description of the bifurcations that occur when one of the  
 287 periodic attractors loses hyperbolicity. The main ingredient in the proof is quasi-symmetric  
 288 rigidity. Recently, Kozlovski announced a simplification of the proof in [21] of this theo-  
 289 rem (using semi-conjugacies to maps with constant absolute value of the slopes, rather than  
 290 stunted sawtooth maps).

### 291 **3.1. Non-local connectivity of isentropes and non-monotonicity in separate variables.**

292 It is possible to parametrize the family  $P^d$  by critical values. The following example shows  
 293 that it is not true that topological entropy depends monotonically on each of these parameters.  
 294 Define  $f_{a,b}(x) = 2ax^3 - 3ax^2 + b$  for  $a = b + 0.515$ . This cubic map has critical points 0  
 295 and 1 and critical values  $f(0) = b$ , and  $f(1) = b - a = 0.515$ . It is shown in [21] that there  
 296 are values of  $b$  such that the map  $a \mapsto h_{top}(f_{a,b})$  is not monotone.

297 Related to this, it is shown in [22] that isentropes in  $P^d$ , when  $d \geq 5$  are not locally  
 298 connected. It is not known whether isentropes in  $P^3$  or in  $P^4$  are locally connected. For  
 299 related results and questions, see [100].

## 300 **4. Measure-theoretical dynamics**

301 We shall now discuss the dynamics of a map  $f : N \rightarrow N$ , where  $N = [0, 1]$  or  $S^1$  from  
 302 measure-theoretical point of view. Recall that a Borel probability measure  $\mu$  is *invariant* for  
 303  $f$  if for each Borel set  $A \subset [0, 1]$  we have  $\mu(f^{-1}A) = \mu(A)$ . We say that  $\mu$  is *ergodic* if a  
 304 Borel set  $A$  with  $f^{-1}(A) = A$  satisfies either  $\mu(A) = 0$  or  $\mu(A) = 1$ . The basin  $B(\mu)$  of  $\mu$   
 305 is the set of points  $x \in [0, 1]$  for which

$$\frac{1}{n} \sum_{i=0}^{n-1} \delta_{f^i(x)} \rightarrow \mu \text{ as } n \rightarrow \infty, \quad (4.1)$$

306 where the convergence is with respect to the weak star topology. If  $B(\mu)$  has positive  
 307 Lebesgue measure, then we say that  $\mu$  is a *physical measure*. Clearly, if  $O$  is an attracting  
 308 periodic orbit, then the averaged Dirac measure  $\mu_O = \frac{1}{\#O} \sum_{p \in O} \delta_p$  is a physical measure.  
 309 An ergodic *acip*, i.e., an invariant probability measure which is absolutely continuous with  
 310 respect to the Lebesgue measure, is also a physical measure, by Birkhoff's ergodic theorem.

311 **4.1. Typical physical measures.** Conjecturally these are the only two types of physical  
 312 measures for typical interval maps, from measure-theoretical point of view. Indeed, in the  
 313 major breakthrough [63], Lyubich proved that *within the quadratic family*  $f_a(x) = ax(1 -$   
 314  $x)$ ,  $1 \leq a \leq 4$ , *for almost every  $a$ , either  $f_a$  is hyperbolic or  $f_a$  has an ergodic acip.* In an  
 315 earlier celebrated work [44], Jakobson showed that the set of  $a$  for which  $f_a$  has an ergodic  
 316 acip has positive Lebesgue measure.

317 An analogue of Lyubich's theorem in the multi-critical case is widely open at the mo-  
 318 ment, due to the multi-dimensional feature of the corresponding parameter space. However,  
 319 a generalization to the case of unimodal polynomials of even degree  $d \geq 2$  is nearly com-  
 320 pleted. The work [6] extends the result of [62], showing that for any even integer  $d \geq 2$ , and  
 321 almost every  $a \in [1, 4]$ ,  $f_a(x) = \frac{a}{4}(1 - (1 - 2x)^d)$  either is hyperbolic, or has an ergodic

322 acip, or is infinitely renormalizable. Moreover, Avila and Lyubich [4] developed a novel way  
 323 to obtain exponential convergence along hybrid classes for infinitely renormalizable maps.  
 324 One can expect a complete proof of the generalization of Lyubich's theorem for unimodal  
 325 maps of a given degree will be available soon. Nevertheless, let us mention in a joint work  
 326 with Bruin, the authors of this paper proved that for all even integer  $d$ , and almost every  
 327  $1 \leq a \leq 4$ ,  $\frac{a}{4}(1 - (1 - 2x)^d)$  has a unique physical measure which might be supported on a  
 328 Cantor set.

329 **4.2. Existence of acip.** We shall now discuss some recent advances on existence of acip for  
 330 smooth interval maps. In order to apply some version of the real Koebe distortion to control  
 331 distortion, we often assume  $f$  lies in the class  $\mathcal{A}_3$  defined below. A map  $f : [0, 1] \rightarrow [0, 1]$   
 332 is in the class  $\mathcal{A}_k$  if the following holds:  $f$  is  $C^1$  and  $C^k$  outside the critical set  $\text{Crit}(f) =$   
 333  $\{c : f'(c) = 0\}$ ; moreover, for each  $c \in \text{Crit}(f)$ , there exists  $\ell_c > 1$  and  $C^k$  diffeomor-  
 334 phisms  $\varphi_c, \psi_c$  of  $\mathbb{R}$  such that  $\varphi_c(c) = \psi_c(f(c)) = 0$  and  $|\psi_c(f(x))| = |\varphi_c(x)|^{\ell_c}$  holds in  
 335 a neighborhood of  $c$ .

336 The following theorem was obtained by the authors in joint with Bruin and Rivera-  
 337 Letelier.

**Theorem 4.1** (Existence of acip [19]). *Let  $f \in \mathcal{A}_3$  be an interval map with all periodic  
 points hyperbolic repelling. Assume that the following large derivatives condition holds: for  
 each  $c \in \text{Crit}(f)$ ,*

$$|Df^n(f(c))| \rightarrow \infty \text{ as } n \rightarrow \infty.$$

338 *Then  $f$  has an acip  $\mu$  with density  $\frac{d\mu}{d\text{Leb}} \in L^p$  for each  $p < \ell_{\max}/(\ell_{\max} - 1)$  where  $\ell_{\max} =$   
 339  $\sup_{c \in \text{Crit}(f)} \ell_c$ .*

340 The unimodal case was done earlier by the authors in joint with Bruin [20]. The existence  
 341 of acip for interval maps has been proved previously in more restrictive settings, including

- 342 • in [72], for maps satisfying the *Misiurewicz* condition:  $\omega(c) \cap \text{Crit}(f) = \emptyset$  for each  
 343  $c \in \text{Crit}(f)$ ;
- 344 • in [27] for unimodal maps satisfying the *Collet-Eckmann* condition (together with other  
 345 conditions): for the critical point  $c$ ,  $\liminf_{n \rightarrow \infty} \frac{1}{n} \log |Df^n(f(c))| > 0$ ;
- 346 • in [74] for unimodal maps satisfying the following summability condition: if  $c$  is the  
 347 critical point and  $\ell$  is the order, then  $\sum_{n=0}^{\infty} |Df^n(f(c))|^{-1/\ell} < \infty$ ,

348 among others. All of the following results assume that  $f$  has negative Schwarzian outside  
 349  $\text{Crit}(f)$  in order to apply the real Koebe principle to control distortion, but now we know that  
 350 the required distortion control is also valid for maps  $f \in \mathcal{A}_3$ , after [48] and [101, Theorem  
 351 C].

352 We should however note that the large derivatives condition is not a necessary condition  
 353 for the existence of an acip, even though an acip necessarily has positive metric entropy:  
 354 there exists a unimodal map in the class  $\mathcal{A}_3$  with  $\liminf |Df^n(f(c))| = 0$  and with an acip  
 355 [16]. It is also known (not surprisingly) that existence of acip is not a topological (or qua-  
 356 siasymmetric) condition [17].

357 **Question 4.2.** *Determine topological (or quasiasymmetric) conjugacy classes in  $\mathcal{A}_3$  such that  
 358 each map in the class has an acip.*

**4.2.1. Ingredients of the proof of Theorem 4.1.** An intermediate step of the proof is to show that the large derivatives condition implies *backward contraction* in the sense of Rivera-Letelier [84], which means the following: if  $\tilde{B}_c(\delta)$  denotes the component of  $f^{-1}(f(c) - \delta, f(c) + \delta)$  which contains  $c$  and

$$\Gamma(\delta) = \inf \left\{ \frac{\delta}{|U|} : \begin{array}{l} U \text{ is a component of } f^{-n}(\tilde{B}_c(\delta)) \text{ containing } f(c') \\ \text{for some } c, c' \in \text{Crit}(f) \text{ and } n \geq 0 \end{array} \right\}$$

359 then  $\Gamma(\delta) \rightarrow \infty$  as  $\delta \rightarrow 0$ . It turns out that the backward contraction property is equivalent  
360 to the large derivatives condition [57].

It is well-known that for any Borel probability measure  $\nu$ , any accumulation point of the following sequence

$$\frac{1}{n} \sum_{i=0}^{n-1} (f^i)_* \nu$$

in the weak star topology is an invariant probability measure of  $f$ , where  $(f^i)_* \nu(A) = \nu(f^{-i}(A))$ . Thus it suffices to prove the following statement: for each  $0 < \kappa < 1$  there exists  $C = C(\kappa)$  such that

$$(f^n)_*(\text{Leb})(A) = |f^{-n}(A)| \leq C |f(A)|^{\kappa/\ell_{\max}},$$

361 holds for all Borel  $A \subset [0, 1]$  and all  $n \geq 0$ . The backward contraction property makes it  
362 possible to obtain the estimate when  $A$  is an interval close to the critical set. For general  $A$ ,  
363 the paper uses a sliding argument from [74], and Mănă's theorem [65].

364 **4.3. Decay of correlation.** A different way to obtain existence of acip is via *inducing*. Let  
365 us say a map  $F : \mathcal{U} \rightarrow \mathcal{V}$ , where  $\mathcal{U} \subset \mathcal{V}$  are open subsets of  $[0, 1]$ , is a *Markov map*, if for  
366 each component  $U$  of  $\mathcal{U}$ ,  $F|_U$  is a  $C^1$  diffeomorphism onto a component of  $\mathcal{V}$ . A Markov  
367 map  $F$  is *induced* by a map  $f$  if there is a continuous function  $s : \mathcal{U} \rightarrow \{1, 2, \dots\}$  such that  
368  $F(x) = f^{s(x)}(x)$ . (So  $s(\cdot)$  takes constant value in each  $U$ .) We shall often consider Markov  
369 maps with extra properties:

- 370 (i)  $\mathcal{V}$  is an interval;
- 371 (i')  $\mathcal{V}$  consists of finitely many intervals;
- (ii) (Bounded distortion) There exist  $C > 0$  and  $\alpha \in (0, 1)$  such that

$$\frac{|DF^n(x)|}{|DF^n(y)|} \leq C |F^n(x) - F^n(y)|^\alpha,$$

372 whenever  $F^i(x)$  and  $F^i(y)$  belong to the same component of  $\mathcal{U}$  for each  $i = 0, 1, \dots$ ,  
373  $n - 1$ .

A Markov map  $F : \mathcal{U} \rightarrow \mathcal{V}$  with the properties (i') and (ii) has an absolutely continuous invariant probability measure  $\nu$  such that  $d\nu/d\text{Leb}$  is bounded away from 0 and  $\infty$ . If we can construct an induced Markov map  $F$  for a map  $f$  such that (i'), (ii) and the following hold:

$$a_s := |\{s(x) \geq s\}| \rightarrow 0 \text{ as } s \rightarrow \infty,$$

then the original system  $f$  has an acip

$$\mu := \frac{1}{\sum_{s=1}^{\infty} a_s} \sum_U \sum_{j=0}^{s|U-1} (f^j)_*(\nu|U),$$

374 where the sum runs over all components of  $\mathcal{U}$ . One advantage of inducing is that through  
375 estimating the speed of convergence of  $a_s \rightarrow 0$ , one can obtain finer statistical properties of  
376 the system.

377 The following theorem was proved by the 1st author in joint with Rivera-Letelier, im-  
378 proving an earlier result [18] considerably.

**Theorem 4.3** (Decay of correlation [85]). *Assume that  $f \in \mathcal{A}_3$  is topologically exact and satisfies the large derivatives condition. Then there is an induced Markov map  $F : \mathcal{U} \rightarrow \mathcal{V}$  such that (i) and (ii) and the following tail estimate hold:*

$$a_s = O(s^{-p}) \text{ for each } p > 0, \text{ as } s \rightarrow \infty.$$

*In particular, the unique acip  $\mu$  of  $f$  is super-polynomially mixing: for each essentially bounded  $\varphi : [0, 1] \rightarrow \mathbb{R}$  and each Hölder continuous  $\psi : [0, 1] \rightarrow \mathbb{R}$ ,*

$$C_n(\varphi, \psi) := \int_0^1 \varphi \circ f^n \psi d\mu - \int_0^1 \varphi d\mu \int_0^1 \psi d\mu$$

379 *converges to 0 superpolynomially fast as  $n \rightarrow \infty$ .*

380 Here we say that  $f$  is topologically exact if for each non-empty open subset  $U$  of  $[0, 1]$ ,  
381 there exists a positive integer  $n$  such that  $f^n(U) = [0, 1]$ . This is a necessary condition  
382 for  $f$  to have a mixing acip. The last statement was deduced from the tail estimate via  
383 Young's tower [102]. Note that the tail estimate also implies finer statistical properties of  
384 the sequence  $\{\psi \circ f^n\}_{n=0}^{\infty}$  (considered as a sequence of random variables with identical dis-  
385 tribution), such as the Central Limit Theorem [102], Almost Sure Invariance Principle [68],  
386 etc, for  $\psi$  Hölder. The paper [85] also dealt with existence and mixing properties of in-  
387 variant probability measures with respect to conformal measures (supported on Julia sets) of  
388 maximal dimension for a large class of complex rational maps. This paper used the induced  
389 Markov map to study the geomtery of the Julia set.

390 Much recent progress on thermodynamical formalism for one-dimensional maps also  
391 used inducing to construct invariant probability measures with respect to various conformal  
392 measures, see for example [23, 76, 82].

For the proof of Theorem 4.3, an adaptation is used of the inducing scheme, called *canonical inducing*, developed in [81, 82]. A crucial new estimate is the following backward shrinking estimate for maps with large derivatives (Theorem B): *there exists  $\rho > 0$  such that*

$$\theta_n := \{|J| : J \text{ is an interval such that } |f^n(J)| \leq \rho\}$$

393 *converges to zero super-polynomially fast.* Theorem C relates the quantity  $\theta_n$  to the tail esti-  
394 mate of a suitably constructed induced Markov map, provided the map has *badness exponent*  
395 0 which was the statement of Theorem A.

396 It is known that  $\theta_n \rightarrow 0$  exponentially fast (the topological Collet-Eckmann condition,  
397 equivalent to the Collet-Eckmann condition in the unimodal case) is equivalent to having an  
398 exponentially mixing acip [73, 81]. It would be interesting to know

399 **Question 4.4.** For a topologically exact interval map  $f \in \mathcal{A}_3$ , is  $\theta_n \rightarrow 0$  superpolynomially  
400 fast equivalent to having a unique acip which is superpolynomially mixing?

401 An affirmative solution to Question 2.11 in [85] implies an affirmative answer to the  
402 question above.

**4.4. Stochastic stability.** An interval map with an acip is not hyperbolic and hence not structurally stable. The notion of stochastic stability, posed by Kolmogorov and Sinai, asks for stability of statistical properties under random perturbations. Given a map  $f : [0, 1] \rightarrow [0, 1]$ , an  $\varepsilon$ -random (pseudo) orbit is by definition a sequence  $\{x_n\}_{n=0}^\infty$  such that  $|f(x_n) - x_{n+1}| \leq \varepsilon$ . Roughly speaking, stochastic stability means when  $\varepsilon > 0$  is small, for most of the  $\varepsilon$ -random orbits  $\{x_n\}_{n=0}^\infty$ , the asymptotic distribution,  $\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \delta_{x_i}$ , is close to a physical measure of  $f$ . Note that if  $f([0, 1]) \subset (0, 1)$  and  $\varepsilon > 0$  small enough, then the space of all  $\varepsilon$ -random orbits can be identified with  $[0, 1] \times [-\varepsilon, \varepsilon]^\mathbb{N}$  by the following formula:

$$\{x_n\}_{n=0}^\infty \mapsto (x_0, x_1 - f(x_0), x_2 - f(x_1), \dots).$$

403 So the space of sequences  $\{x_n\}_{n=0}^\infty$  can be endowed with a probability measure  $\mathbb{P}_\varepsilon$  which  
404 corresponds to  $m \times m_\varepsilon^\mathbb{N}$ , where  $m$  denotes the Lebesgue measure on  $[0, 1]$  and  $m_\varepsilon$  denotes  
405 the normalised Lebesgue measure on  $[-\varepsilon, \varepsilon]$ . In the literature, reference measures other than  
406  $\mathbb{P}_\varepsilon$  have also been considered on the space of  $\varepsilon$ -random orbits, corresponding to different  
407 types of random perturbations. The measure  $\mathbb{P}_\varepsilon$  corresponds to the so-called *additive noise*  
408 model.

409 Recently the 1st author proved the following theorem.

410 **Theorem 4.5** (Stochastic Stability [88]). Suppose  $f \in \mathcal{A}_3$  is ergodic with respect to the  
411 Lebesgue measure and that the following summability condition holds: for each  $c \in \text{Crit}(f)$ ,

$$\sum_{n=0}^{\infty} |Df^n(f(c))|^{-1} < \infty. \quad (4.2)$$

Then the unique acip of  $f$  is stochastic stable in the strong sense: For each  $\varepsilon > 0$  there exists a unique probability measure  $\mu_\varepsilon$  absolutely continuous with respect to the Lebesgue measure, such that for  $\mathbb{P}_\varepsilon$ -a.e.  $\varepsilon$ -random orbits  $\{x_n\}_{n=0}^\infty$ ,

$$\frac{1}{n} \sum_{i=0}^{n-1} \delta_{x_i} \rightarrow \mu_\varepsilon$$

412 as  $n \rightarrow \infty$  in the weak star topology. Moreover, the density  $\frac{d\mu_\varepsilon}{d\text{Leb}}$  converges in  $L^1$  to the  
413 density of the unique acip of  $f$  as  $\varepsilon \rightarrow 0$ .

414 See the Main Theorem of [88] for a more general statement, which covers a very general  
415 type of random perturbation. Previously, stochastic stability was studied for interval maps  
416 with a Benedicks-Carleson type condition [12, 13] (or even stronger) which thus has expo-  
417 nential decay of correlation, see [11, 14, 46, 95]. It is surprising that the stochastic stability  
418 of the Manneville-Pomeau map  $x \mapsto x + x^{1+\alpha} \pmod 1$ , which is probably the simplest non-  
419 uniformly expanding dynamical system, was only established very recently by the authors  
420 in [89].

Li and Wang [59] proved stochastic stability for unimodal maps  $f$  with a wild attractor where the physical measure is supported on the Cantor attractor. It raises a curious question whether there exists an interval map with a stochastically unstable physical measure.

The crucial step in the proof of Theorem 4.5 was to establish the first estimate on the first return maps to critical neighborhoods: Let  $\varepsilon > 0$  be small and let  $\tilde{B}_c(\varepsilon)$  be defined as in § 4.2.1. Then for all  $\varepsilon$ -random orbits  $\{x_i\}_{i=0}^n$  with  $x_0 \in \tilde{B}_{c_1}(\varepsilon)$ ,  $x_n \in \tilde{B}_{c_2}(\varepsilon)$  for some  $c_1, c_2 \in \text{Crit}(f)$  and  $x_1, x_2, \dots, x_{n-1} \notin \bigcup_{c \in \text{Crit}(f)} \tilde{B}_c$ , we have

$$\prod_{i=1}^{n-1} |Df(x_i)| \geq \frac{\Lambda(\varepsilon)}{\varepsilon^{1-\ell_{c_2}^{-1}}} \exp(\varepsilon^{\alpha(\varepsilon)} n),$$

where  $\Lambda(\varepsilon) \rightarrow \infty$  and  $\alpha(\varepsilon) \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . The measure  $\mu_\varepsilon$  was constructed using a random inducing scheme initiated in [8]. See also [1, 2].

**4.5. Jakobson’s theorem.** The lower bound for derivative plays a crucial role in a generalization of Jakobson’s theorem by B. Gao and the 1st author [38]. Among a huge number of works in generalizing Jakobson’s theorem, our approach is close to that of [96]. While the paper worked with general one-parameter families, the following is the main result obtained for polynomial maps.

**Theorem 4.6** (Summability implies Collet-Eckmann almost surely [38]). *Fix an integer  $n \geq 2$ . For each  $\mathbf{a} = (a_0, a_1, \dots, a_n) \in \mathbb{R}^{n+1}$  write  $P_{\mathbf{a}}(x) = \sum_{i=0}^n a_i x^i$ . Let  $\Lambda_n$  denote the collection of  $\mathbf{a} \in \mathbb{R}^{n+1} \setminus \{\mathbf{0}\}$  for which the following hold: (i)  $P_{\mathbf{a}}([0, 1]) \subset [0, 1]$  and (ii)  $P_{\mathbf{a}} : [0, 1] \rightarrow [0, 1]$  satisfies the summability condition (4.2). Then  $\Lambda_n$  has positive measures and almost every  $\mathbf{a} \in \Lambda_n$  satisfies the Collet-Eckmann condition, and the following polynomial recurrence conditions: for each  $\beta > 1$ , and any critical points  $c, c'$  of  $P_{\mathbf{a}}|_{[0, 1]}$ , we have  $|P_{\mathbf{a}}^k(c) - c'| \geq k^{-\beta}$  for all  $k$  sufficiently large.*

The proof is done by purely real analytic method, except we had to use a recent transversality result due to Levin [54] which was based on complex methods. For the case  $n = 2$ , the transversality result was known before in [3, 53]. For the quadratic family, the Collet-Eckmann and polynomial recurrence conditions are satisfied by almost every non-hyperbolic map [7]. It would be interesting to push the real analytic method further, for instance, to see whether the summability condition can be replaced by the large derivatives condition in Theorems 4.5 and 4.6.

Finally we would like to draw the reader’s attention to the works [9, 10] where the “modulus of continuity” of  $t \mapsto \mu_t$  over “good” non-uniformly expanding maps is studied for families  $f_t$  of unimodal maps, where  $\mu_t$  is the acip for  $f_t$ .

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National University of Singapore and Imperial College, London

692 E-mail: matsw@nus.edu.sg; s.van-strien@imperial.ac.uk