# Math 273 Lecture 21 

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At the beginning of the course, we used open book decompositions to show that every closed 3 -manifold admits a contact structure. There is a strong relationship between these two notions, due to Giroux [1], which we will now investigate; first, we recall some definitions as well as Thurston and Winkelnkemper's proof that contact structures exist.

Definition 1. An open book decomposition of a closed 3-manifold $Y$ is a pair $\left(B, \pi: Y \backslash B \rightarrow S^{1}\right)$, where $B \subset Y$ is an oriented link called the binding and $\pi$ is a fibration. The fibers $\pi^{-1}(\theta)$ are the interiors of compact surfaces $\Sigma_{\theta}$ called pages, which satisfy $B=\partial \Sigma_{\theta}$ for all $\theta$.

Theorem 2 ([3]). An open book decomposition $(B, \pi)$ gives rise to a contact structure on $Y$.

Proof (sketch). Let $\Sigma$ be the page of the open book, and write $Y \backslash N(B)$ as the mapping torus of $h: \Sigma \rightarrow \Sigma$. Take a 1-form $\lambda$ on $\Sigma$ such that $d \lambda$ is an area form and $\lambda=(1+t) d \theta$ on a neighborhood $[-1,0]_{t} \times S_{\theta}^{1}$ of each component of $\partial \Sigma$; the space of all such $\lambda$ is convex and nonempty. Then the 1 -form

$$
\alpha=\phi \lambda+(1-\phi) h^{*} \lambda+K d \phi
$$

is a contact form on $\Sigma \times[0,1]_{\phi}$ for $K$ large, and it descends to a contact form on the mapping torus $Y \backslash N(B)$. Near $\partial N(B)$ it is equal to $(1+t) d \theta+K d \phi$.

Now on each solid torus component $S_{\theta}^{1} \times D_{(r, \phi)}^{2}$ of $N(B)$, we need to identify a contact form $\alpha$ which is equal to $(2-r) d \theta+K d \phi$ near $r=1$; we will take $\alpha=f(r) d \theta+g(r) d \phi$ for some functions $f$ and $g$ which equal $2-r$ and $K$ respectively near $r=1$. We pick $f$ and $g$ to satisfy $f(r)=1$ and $g(r)=r^{2}$ near $r=0$ and then choose any extension of $(f, g)$ over the rest of $[0,1]$ which satisfies $f g^{\prime}-g f^{\prime}>0$; this is equivalent to the contact condition $\alpha \wedge d \alpha>0$, so we are done.

Near any component of the binding $B$ we have $\alpha=d \theta+r^{2} d \phi$, where $B$ given in coordinates by $r=0$ and $\partial_{\theta}$ is tangent to $B$ with the same orientation it inherits as $\partial \Sigma_{\theta}$. Since $\alpha\left(\partial_{\theta}\right)=1, B$ is a positively transverse link. Furthermore, on any page $\Sigma_{\theta}$ we have $\left.d \alpha\right|_{\Sigma_{\theta}}=\phi d \lambda+(1-\phi) h^{*} d \lambda$, which is an area form because both $d \lambda$ and $h^{*} d \lambda$ are. We describe these properties as follows:

Definition 3. A contact structure $\xi$ on $Y$ is supported by the open book $(B, \pi)$ if up to isotopy it admits a contact form $\alpha$ such that $B$ is a positively transverse link and $d \alpha$ is a positive area form on each page $\Sigma_{\theta}$.

Thus Theorem 2 actually provides a contact structure supported by the given open book.

Example 4. Consider $S^{3}$ as the unit sphere in $\mathbb{C}^{2}$, with the standard contact form $\alpha=r_{1}^{2} d \theta_{1}+r_{2}^{2} d \theta_{2}$. The positive Hopf link $H \subset S^{3}$ can be expressed as the zero set of the polynomial $f\left(z_{1}, z_{2}\right)=z_{1} z_{2}$, and then the map

$$
\pi\left(z_{1}, z_{2}\right)=\frac{z_{1} z_{2}}{\left|z_{1} z_{2}\right|}
$$

defines a fibration $\pi: S^{3} \backslash H \rightarrow S^{1}$. Each page $\Sigma_{\phi}=\pi^{-1}(\phi)$ can be parametrized by a map

$$
f(r, \theta)=(\sqrt{t}, \theta, \sqrt{1-t}, \phi-\theta)
$$

in polar coordinates, i.e. $z_{1}=\sqrt{t} e^{i \theta}$ and $z_{2}=\sqrt{1-t} e^{i(\phi-\theta)}$; here $0<t<1$ and $\theta \in S^{1}$, so the pages are annuli. Then

$$
f^{*} \alpha=t d \theta+(1-t)(-d \theta)=(2 t-1) d \theta
$$

and so $f^{*} d \alpha=2 d t \wedge d \theta$, which is an area form on $[0,1]_{t} \times S_{\theta}^{1}$, so $d \alpha$ is an area form on each page. Furthermore, the components $\left\{r_{1}=0\right\}$ and $\left\{r_{2}=0\right\}$ of $H$ have oriented tangent vectors $\partial_{\theta_{2}}$ and $\partial_{\theta_{1}}$ respectively, at which points $\alpha\left(\partial_{\theta_{2}}\right)=r_{2}^{2} \iota_{\partial_{\theta_{2}}} d \theta_{2}=1$ and $\alpha\left(\partial_{\theta_{1}}\right)=r_{1}^{2} \iota \iota_{\theta_{1}} d \theta_{1}=1$ are positive. Thus $(H, \pi)$ supports the standard tight contact structure $\xi_{\mathrm{st}}$ on $S^{3}$.

Exercise 5. Show that the open book decomposition of $S^{3}$ whose binding is the unknot $U=\left\{r_{1}=0\right\}$ and whose pages are the disks $\left\{\theta_{1}=\right.$ const. $\}$ also supports $\xi_{\text {st }}$.

Proposition 6. Let $\xi$ and $\xi^{\prime}$ be contact structures supported by the same open book $(B, \pi)$. Then $\xi$ is isotopic to $\xi^{\prime}$.

Proof. We will first construct isotopies of these contact structures so that they have sufficiently nice contact forms. Apply an isotopy so that $\xi$ has a contact form $\alpha$ with the desired properties. Then we let $d \phi$ be an area form on $S^{1}$ and pull it back to get a form $\pi^{*} d \phi$ on $Y \backslash B$. In a neighborhood $N=S^{1} \times D^{2}$ of any component $B_{i}$ of the binding, we can choose coordinates $(\theta,(r, \phi))$ so that $d \phi$ agrees with the form $\pi^{*} d \phi$, and since $\alpha\left(\partial_{\theta}\right)>0$ along $B_{i}$ we can also insist that $\alpha\left(\partial_{\theta}\right)>0$ on all of $N$ by possibly shrinking it. Take $\epsilon>0$ small enough that the solid torus $N^{\prime}=\{r<\epsilon\}$ lies inside $N$, and define a smooth nondecreasing function $f:[0, \epsilon] \rightarrow[0,1]$ so that $f(r)=r^{2}$ near $r=0$ and $f(r)=1$ near $r=\epsilon$. Then $f$ defines a function on $N^{\prime}$ which we can extend to all of $Y$ by letting $f=1$ outside any of the solid tori $N^{\prime}$.

For any $t \geq 0$, we define a new 1-form $\alpha_{t}=\alpha+t f(r) d \phi$. We can check that

$$
\alpha_{t} \wedge d \alpha_{t}=\alpha \wedge d \alpha+t f(r) d \phi \wedge d \alpha+t \alpha \wedge d f \wedge d \phi
$$

Note that $t f(r) d \phi \wedge d \alpha \geq 0$, because $d \alpha$ is a volume form on each page and $d \phi$ is dual to a vector which is positively transverse to each page. Furthermore, $\alpha \wedge d f \wedge d \phi$ is zero outside the neighborhoods $N^{\prime}$ (where $d f=0$ ) and equal to $t f^{\prime}(r) \cdot \alpha \wedge d r \wedge d \phi$ on $N^{\prime}$; since $d r \wedge d \phi$ is nonnegative on $D^{2}$ and zero on $\partial_{\theta}$, while $\alpha\left(\partial_{\theta}\right)>0$, we conclude that $t \alpha \wedge d f \wedge d \phi \geq 0$ as well. Therefore $\alpha_{t}$ is actually a contact form for any $t>0$, and since $\alpha=\alpha_{0}$ these all define contact structures isotopic to $\xi$. We use the same construction, with the same neighborhoods $N^{\prime}$ of $B$, to find contact forms $\alpha_{t}^{\prime}$ for $\xi^{\prime}$ up to isotopy.

Finally, we take $t$ to be very large and define an isotopy from $\xi$ to $\xi^{\prime}$ by the family of contact forms $\alpha_{s}=(1-s) \alpha_{t}+s \alpha_{t}^{\prime}$. Then we compute

$$
\alpha_{s} \wedge d \alpha_{s}=\left[(1-s)^{2} \alpha_{t} \wedge d \alpha_{t}+s^{2} \alpha_{t}^{\prime} \wedge d \alpha_{t}^{\prime}\right]+s(1-s)\left[\alpha_{t} \wedge d \alpha_{t}^{\prime}+\alpha_{t}^{\prime} \wedge d \alpha_{t}\right]
$$

and the first term in brackets is clearly positive, so we must show that the second term in brackets is nonnegative. Outside the neighborhoods $N^{\prime}$, where $f=1$ and $d f=0$, we have

$$
\alpha_{t} \wedge d \alpha_{t}^{\prime}+\alpha_{t}^{\prime} \wedge d \alpha_{t}=(\alpha+t d \phi) \wedge d \alpha^{\prime}+\left(\alpha^{\prime}+t d \phi\right) \wedge d \alpha=t d \phi \wedge\left(d \alpha+d \alpha^{\prime}\right)+O(1)
$$

which is positive for $t$ large enough. Inside the neighborhoods $N^{\prime}$, we have

$$
\begin{aligned}
\alpha_{t} \wedge d \alpha_{t}^{\prime}+\alpha_{t}^{\prime} \wedge d \alpha_{t} & =(\alpha+t f d \phi) \wedge\left(d \alpha^{\prime}+t f^{\prime} d r \wedge d \phi\right)+\left(\alpha^{\prime}+t f d \phi\right) \wedge\left(d \alpha+t f^{\prime} d r \wedge d \phi\right) \\
& =t\left(f^{\prime} \cdot\left(\alpha+\alpha^{\prime}\right) \wedge d r \wedge d \phi+f d \phi \wedge\left(d \alpha+d \alpha^{\prime}\right)\right)+O(1)
\end{aligned}
$$

Note that $f^{\prime}$ and $f$ are both nonnegative. Near $r=0$ we have $f^{\prime} d r \wedge d \phi=$ $2 r d r \wedge d \theta$, which is a positive area form on $D^{2}$ and so the first term in parentheses is positive; then away from $r=0$ we have $f>0$ and so the second term is positive. Thus for $t$ large enough this form is positive on all of $N^{\prime}$, and so $\alpha_{s}$ is a contact form for $0 \leq s \leq 1$. We conclude that it gives the desired isotopy between $\xi$ and $\xi^{\prime}$.

We have now shown that every open book decomposition of $Y$ supports a unique contact structure up to isotopy. The Giroux correspondence provides a converse: every contact structure is supported by an open book, and furthermore if two open books support the same contact structure then they are related by a series of moves called positive stabilizations. Before proving this, we will consider some equivalent ways to describe the contact structure supported by an open book.

Proposition 7. The contact structure $(Y, \xi)$ is supported by the open book $(B, \pi)$ if and only if up to isotopy, $\xi$ admits a contact form whose Reeb vector field $R$ is positively tangent to $B$ and positively transverse to the pages of $\pi$.

Proof. Suppose we have a contact form $\alpha$ with such a Reeb vector field $R$, and recall that $\alpha(R)=1$ and $\iota_{R} d \alpha=0$ by definition. Then $\alpha(R)>0$ along $B$, so $B$ is a positively transverse link. At any point $p$ on a page $\Sigma_{\theta}$, we have $\iota_{R}(\alpha \wedge d \alpha)=\left(\iota_{R} \alpha\right) d \alpha-\alpha \wedge \iota_{R} d \alpha=d \alpha$, and so given a positive basis $\left(v_{1}, v_{2}\right)$ for $T_{p} \Sigma_{\theta}$ we have

$$
d \alpha\left(v_{1}, v_{2}\right)=(\alpha \wedge d \alpha)\left(R, v_{1}, v_{2}\right)>0
$$

so $d \alpha$ is an area form on $\Sigma_{\theta}$. Therefore $(B, \pi)$ supports $\xi=\operatorname{ker}(\alpha)$.
Conversely, suppose that $(Y, \xi)$ is supported by $(B, \pi)$, so that there is a contact form $\alpha$ for which $B$ is positively transverse and $d \alpha$ is positive on each page. Let $R$ be the Reeb vector field of $\alpha$. Then at any point on a page, the fact that $\iota_{R}(\alpha \wedge d \alpha)=d \alpha$ is an area form on the page implies that $R$ must be positively transverse to that page. Furthermore, on a neighborhood $S_{\theta}^{1} \times D^{2}$ of a component of $B$, we can write $R=a \partial_{\theta}+b \partial_{x}+c \partial_{y}$ where $(x, y)$ are rectangular coordinates on $D^{2}$. If at $(x, y)=(0,0)$ we have $(b, c) \neq(0,0)$, then either $b$ or $c$ is nonzero on an entire neighborhood of the point $(\theta, 0,0)$. The oriented normal vectors to the pages $\Sigma_{0}, \Sigma_{\pi / 2}, \Sigma_{\pi}$, and $\Sigma_{3 \pi / 2}$ at points a distance $\epsilon$ away from $B$ are given by $\partial_{y},-\partial_{x},-\partial_{y}$, and $\partial_{x}$ respectively, and so these pages cannot be positively transverse to $R$ in the respective cases $c<0, b>0, c>0$, and $b<0$. We conclude that in fact $(b, c)=(0,0)$ along $B$, so the Reeb vector has the form $R=a \partial_{\theta}$, and then $a \cdot \alpha\left(\partial_{\theta}\right)=\alpha(R)=1$ and $\alpha\left(\partial_{\theta}\right)>0$ together imply that $a>0$, so $R$ is positively tangent to $B$.

Definition 8. Let $f: \mathbb{C}^{2} \rightarrow \mathbb{C}$ be a polynomial with $f(0,0)=0$ and an isolated singularity at the origin. Then for $\epsilon>0$ small, the sphere $S^{3}$ of radius $\epsilon$ centered at the origin admits a Milnor open book $(B, \pi)$, where $B=f^{-1}(0) \cap S^{3}$ and $\pi: S^{3} \backslash B \rightarrow S^{1}$ is the Milnor fibration

$$
\pi\left(z_{1}, z_{2}\right)=\frac{f\left(z_{1}, z_{2}\right)}{\left|f\left(z_{1}, z_{2}\right)\right|}
$$

Proposition 9. A Milnor open book on $S^{3}$ supports the tight contact structure.
Proof. We will give a proof in the case where $f$ is homogeneous. Recall that the standard contact structure $\xi_{\text {st }}$ is given by $T S^{3} \cap i\left(T S^{3}\right)$ corresponding to the strictly plurisubharmonic exhaustion function $\phi\left(z_{1}, z_{2}\right)=\frac{1}{2}\left(\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}\right)$ on $B^{4}$.

We wish to show that the Reeb vector field $R$ is positively transverse to the pages of $\pi$, so first we must determine $R$; we will compute $R$ for the boundary of a general Stein domain and then specialize to the case $S^{3}=\partial B^{4}$. Given an spsh function $\phi: X \rightarrow \mathbb{R}$ and complex structure $J$, the induced contact form on the level set $\partial X$ is $\alpha=-d \phi \circ J$. Write the Reeb vector field in the form $R=J v$ for some vector $v$; then we have an induced metric $\langle x, y\rangle=d \alpha(x, J y)$ on $X$, so $\left(\iota_{R} d \alpha\right)(x)=d \alpha(J v, x)=-\langle x, v\rangle$. Since $\iota_{R} d \alpha=0$ on $T S^{3}$, the vector $v$ must be orthogonal to all of $T(\partial X)$, hence it is a multiple of $\nabla \phi$. Furthermore, $R$ satisfies $\alpha(R)=-d \phi(J R)=1$, so $\left\langle\nabla \phi,-J^{2} v\right\rangle=-1$ or simply $\langle\nabla \phi, v\rangle=1$. Therefore $R=J \frac{\nabla \phi}{|\nabla \phi|^{2}}$.

Next, we observe that

$$
\nabla \phi=\sum_{i=1}^{2}\left(x_{i} \partial_{x_{i}}+y_{i} \partial_{y_{i}}\right)
$$

so since $f$ is homogeneous it follows that $d f(\nabla \phi)=(\operatorname{deg} f) f$. In particular, along the level set $B=f^{-1}(0)$ we have $d f(R)=d f\left(\frac{i}{|\nabla \phi|^{2}} \nabla \phi\right)=0$ because
$d f(\nabla \phi)=(\operatorname{deg} f) \cdot 0=0$, and so $R$ is tangent to $B$. Next, by writing $f$ in polar coordinates and taking logarithmic derivatives, we see that $d \pi=\operatorname{Im}\left(f^{-1} d f\right)$. We want to show that $R$ is transverse to the pages of $\pi$, so it will suffice to show that $d \pi(R)>0$, or equivalently that $d \pi(i \nabla \phi)>0$ since the two differ by a factor of $|\nabla \phi|^{2}$. Now we can compute

$$
d \pi(i \nabla \phi)=\operatorname{Im}\left(f^{-1} d f(i \nabla \phi)\right)=\operatorname{Re}\left(f^{-1} d f(\nabla \phi)\right)=\operatorname{Re}(\operatorname{deg} f)=\operatorname{deg} f>0
$$

as desired. Therefore $(B, \pi)$ supports $\xi_{\text {st }}$.
Now suppose $f$ is not homogeneous. We replace $\phi\left(z_{1}, z_{2}\right)=\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}$ with the function

$$
\psi\left(z_{1}, z_{2}\right)=\frac{1}{2}\left(\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\frac{1}{c}\left|f\left(z_{1}, z_{2}\right)\right|^{2}\right)
$$

and let $S^{3}$ be the level set $\psi^{-1}\left(\frac{r^{2}}{2}\right)$, where $c$ and $r$ are some positive constants. The Reeb vector field on $S^{3}$ is now $R=i \frac{\nabla \psi}{|\nabla \psi|^{2}}$, where

$$
\nabla \psi=\nabla \phi+\frac{1}{c} f \nabla f
$$

and since $d \pi=\operatorname{Im}\left(f^{-1} d f\right)$ as before we have

$$
d \pi(R)=\frac{1}{|\nabla \psi|^{2}} \operatorname{Re}\left(f^{-1} d f\left(\nabla \phi+\frac{1}{c} f \nabla f\right)\right)=\frac{1}{|\nabla \psi|^{2}} \operatorname{Re}\left(\frac{d f(\nabla \phi)}{f}+\frac{d f(\nabla f)}{c}\right) .
$$

Since $d f(\nabla f)=|\nabla f|^{2}$ and $\frac{1}{c}=\frac{r^{2}-|z|^{2}}{|f|^{2}}$ along $S^{3}=\psi^{-1}\left(\frac{r^{2}}{2}\right)$, we conclude that up to a factor of $|\nabla \psi|^{2}, d \pi(R)$ is the real part of

$$
\frac{d f(\nabla \phi)}{f}+\left(r^{2}-|z|^{2}\right)\left|\frac{\nabla f}{f}\right|^{2} .
$$

We leave it as an exercise to show that this quantity can also be made positive on $S^{3} \backslash f^{-1}(0)$ by taking $r$ sufficiently small and $c$ sufficiently large compared to $r$.

## References

[1] Emmanuel Giroux, Contact structures and symplectic fibrations over the circle, notes from the summer school "Holomorphic curves and contact topology, Berder, June 2003", available at http://web.archive.org/web/20040828230626/http://www-fourier.ujfgrenoble.fr/~eferrand/giroux3.ps.
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[3] W. P. Thurston and H. E. Winkelnkemper, On the existence of contact forms, Proc. Amer. Math. Soc. 52 (1975), 345-347.

