

COMPLEX AND REAL VARIABLE HELE-SHAW MOVING BOUNDARY PROBLEM

Sergei ROGOSIN^{1,2}, Gennady MISHURIS²

¹Belarusian State University, 4, Nezavisimosti ave, 220030 Minsk, Belarus

²Aberystwyth University, Penglais, SY23 3BZ, Aberystwyth, UK

Email: rogosinsv@gmail.com; ggm@aber.ac.uk

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The aims:

- to report on the method of reduction of the Hele-Shaw model to an abstract Cauchy-Kovalevsky problem;
- to present results for unbounded and bounded Hele-Shaw cell;
- to give a formulation of the real-variable Hele-Shaw model;
- to describe the way of application of certain asymptotic methods.

Plan:

- ◇ Complex-variable Hele-Shaw model
- ◇ Mathematical model for Hele-Shaw problem
- ◇ Nirenberg-Nishida-Ovsjannikov theorem
- ◇ Case 1. Complex-variable Hele-Shaw model with kinetic undercooling regularization
- ◇ Case 2. Complex-variable Hele-Shaw flow in a bounded cell.
- ◇ Real-variable Hele-Shaw model

Hele-Shaw model. *Hele-Shaw flow* is the flow of a viscous fluid between two closely related parallel plate. The Navier-Stokes equations neglecting the gravity

$$\frac{\partial \mathbf{V}}{\partial t} + (\mathbf{V} \cdot \nabla) \mathbf{V} = \frac{1}{\rho} (-\nabla p + \mu \Delta \mathbf{V}), \quad \nabla \cdot \mathbf{V} = 0. \quad (1)$$

Assumptions

$$\frac{\partial \mathbf{V}}{\partial t} = 0, \quad V_3 = 0.$$

Under these the equations (1) become

$$\begin{aligned} \left(V_1 \frac{\partial}{\partial x_1} + V_2 \frac{\partial}{\partial x_2} \right) V_1 &= -\frac{1}{\rho} \frac{\partial p}{\partial x_1} + \frac{\mu}{\rho} \Delta V_1, \\ \left(V_1 \frac{\partial}{\partial x_1} + V_2 \frac{\partial}{\partial x_2} \right) V_2 &= -\frac{1}{\rho} \frac{\partial p}{\partial x_2} + \frac{\mu}{\rho} \Delta V_2, \\ 0 &= -\frac{1}{\rho} \frac{\partial p}{\partial x_3}, \end{aligned}$$

with boundary conditions

$$V_1 \Big|_{x_3=0,h} = V_2 \Big|_{x_3=0,h} = 0.$$

If h is sufficiently small and the flow is slow, then

$$\frac{\partial p}{\partial x_1} = \mu \frac{\partial^2 V_1}{\partial x_3^2}, \quad \frac{\partial p}{\partial x_2} = \mu \frac{\partial^2 V_2}{\partial x_3^2}, \quad 0 = \frac{\partial p}{\partial x_3}.$$

The boundary conditions then imply

$$V_1 = \frac{1}{2} \frac{\partial p}{\partial x_1} \left(\frac{x_3^2}{\mu} - \frac{hx_3}{\mu} \right), \quad V_2 = \frac{1}{2} \frac{\partial p}{\partial x_2} \left(\frac{x_3^2}{\mu} - \frac{hx_3}{\mu} \right).$$

The integral mean $\widetilde{\mathbf{V}}$ of \mathbf{V} satisfies *Hele-Shaw equation*

$$\widetilde{\mathbf{V}} = -\frac{h^2}{12\mu} \nabla p. \tag{2}$$

A point sink/source (x_1^0, x_2^0) of constant strength. Then

$$\iint_{U_\varepsilon} \frac{h^2 \rho}{12\mu} \Delta p dx_1 dx_2 = \text{const.}$$

On the fluid boundary the balance of forces gives

$$p = \text{exterior air pressure} + \text{surface tension.}$$

Let $\Omega(t)$ be the bounded simply connected domain and consider suction/injection through a single sink/source at the origin. $\Gamma(t) = \partial\Omega(t)$ ($\Omega(0) =: \Omega_0, \Gamma(0) =: \Gamma_0$).

$$\Delta p = Q\delta_0(z), \quad z = x + iy \in \Omega(t). \quad (3)$$

The zero surface tension *dynamic* boundary condition is given by

$$p(z, t) = 0, \quad \forall z \in \Gamma(t). \quad (4)$$

$$v_n = \mathbf{V} \Big|_{\Gamma(t)} \cdot \mathbf{n}(t), \quad (5)$$

The *kinematic* boundary condition

$$\frac{\partial p}{\partial \mathbf{n}} = -v_n. \quad (6)$$

The complex potential $\chi(z, t)$, $\text{Re } \chi = p$.

$$\frac{\partial \chi}{\partial z} = \frac{\partial p}{\partial x_1} - i \frac{\partial p}{\partial x_2}.$$

$$\chi(z, t) = \frac{Q}{2\pi} \log z + \chi_0(z, t). \quad (7)$$

The Polubarinova-Galin equation: find a family of conformal mappings $f(z, t) : G_1 = \{z \in \mathbb{C} : |z| < 1\} \rightarrow \Omega(t)$, s.t.

$$\text{Re} \left\{ f'(\zeta, t) \overline{\zeta f'(\zeta, t)} \right\} = -\frac{Q}{2\pi}, \quad \zeta = e^{i\phi}. \quad (8)$$

Mathematical model for the complex-variable Hele-Shaw problem.

$$\text{Re} \left(\frac{1}{z} \frac{\partial f}{\partial t}(z, t) \overline{\frac{\partial f}{\partial z}(z, t)} \right) = -\frac{Q(t)}{2\pi}, \quad (z, t) \in \partial G_1 \times [0, T]; \quad (9)$$

$$f(z, 0) = f_0(z), \quad z \in G_1; \quad (10)$$

$$f(0, t) = 0, \quad t \in [0, T]. \quad (11)$$

Nirenberg-Nishida-Ovsjannikov theorem. *Let us consider the problem*

$$d_t w = F(t, w), \quad w(0) = 0, \quad (12)$$

in a scale of Banach spaces $\{B_s, \|\cdot\|_s\}_{0 < s \leq 1}$, and let

- *for each fixed $s, 0 < s \leq 1$, the mapping $F(t, w)$ of $[0, T] \times \{w \in B_s : \|w\|_s < R\}$ to $B_{s'}, 0 < s' < s$, is continuous with respect to t ;*
- *for all $0 < s' < 1$ the continuous function $F(t, 0)$ satisfies*

$$\|F(t, 0)\|_{s'} \leq \frac{K}{1 - s'};$$

- *for all $0 < s' < s \leq 1, t \in [0, T], w_1, w_2 \in \{\|w\|_s < R\}$*

$$\|F(t, w_1) - F(t, w_2)\|_{s'} \leq \frac{C}{s - s'} \|w_1 - w_2\|_s.$$

Then the problem (12) has a unique solution

$$w \in \mathcal{C}^1([0, a_0(1 - s)), B_s)_{0 < s < 1}, \quad \|w(t)\|_s < R,$$

where a_0 is a suitable positive constant.

Case 1. Complex-variable Hele-Shaw model with kinetic under-cooling regularization. Problem (P_α):

$$\operatorname{Re} \left(\frac{1}{z} \frac{\partial f}{\partial t} (z, t) \overline{\frac{\partial f}{\partial z} (z, t)} \right) = -\frac{Q(t)}{2\pi} + \operatorname{Re} (zw_{reg}(z, t)), \quad (13)$$

$$f(z, 0) = f_0(z), \quad (14)$$

$$f(0, t) = 0, \quad (15)$$

$$\operatorname{Im} (zw_{reg}(z, t)) = \alpha \partial_\theta \left(\left| \frac{\partial f}{\partial z} \right|^{-1} \left(Q(t) + \operatorname{Re} (zw_{reg}(z, t)) \right) \right), \quad (16)$$

where $\alpha > 0, z = e^{i\theta}$.

For complex potential $\chi = \chi(f(z, t), t)$ we have

$$\operatorname{Re} \left(\frac{1}{z} \frac{\partial f}{\partial t} \overline{\frac{\partial f}{\partial z}} \right) = \operatorname{Re} \left(z \frac{\partial \chi}{\partial z} \right),$$

$$\operatorname{Re} \chi = -\alpha \left| \frac{\partial f}{\partial z} \right|^{-1} \operatorname{Re} \left(\frac{1}{z} \frac{\partial f}{\partial t} \overline{\frac{\partial f}{\partial z}} \right), \quad \alpha > 0.$$

The ansatz for $\chi = Q(t) \log z + \chi_{reg}$, and for $\bar{w} := \partial_z \chi = Q(t)/z + w_{reg}$.

On structure of the problem (P_α). Under the additional assumption

$$\operatorname{Im} \left(\frac{1}{z} \frac{\partial f}{\partial t}(z, t) \left(\frac{\partial f}{\partial z} \right)^{-1}(z, t) \right) (0, t) = 0, \quad (17)$$

one can rewrite (13) by using Schwarz's integral formula as follows:

$$\frac{\partial f}{\partial t}(z, t) - z \frac{\partial f}{\partial z}(z, t) \frac{1}{2\pi i} \int_{|\zeta|=1} \left| \frac{\partial f}{\partial \zeta} \right|^{-2} (Q(t) + \operatorname{Re}(\zeta w_{reg})) \frac{\zeta + z d\zeta}{\zeta - z \zeta} = 0.$$

Differentiating and setting $\omega(z, t) := z w_{reg}(z, t)$, $\phi(z, t) := \left(\frac{\partial f}{\partial z}\right)^{-1}(z, t)$ gives

Problem (Q_α):

- abstract Cauchy-Kovalevsky problem

$$\frac{\partial \phi}{\partial t}(z, t) - z \frac{\partial \phi}{\partial z}(z, t) \mathbf{T}_t(\phi, \omega) + \phi(z, t) \frac{\partial}{\partial z} (z \mathbf{T}_t(\phi, \omega)) = 0, \quad (18)$$

$$\phi_0(z) := \phi(z, 0) = \left(\frac{\partial f_0}{\partial z}\right)^{-1}, \quad \mathbf{T}_t(\phi, \omega) := \frac{1}{2\pi i} \int_{|\zeta|=1} |\phi|^2 (Q(t) + \operatorname{Re} \omega) \frac{\zeta + z d\zeta}{\zeta - z \zeta}. \quad (19)$$

- Riemann-Hilbert-Poincaré problem for $\omega = \omega(z, t)$:

$$\operatorname{Im} (\omega(z, t)) = \alpha \partial_\theta \left(|\phi(z, t)| \left(Q(t) + \operatorname{Re} (\omega(z, t)) \right) \right) \text{ on } \partial G_1 \times [0, T]. \quad (20)$$

Mathematical treatment of problem (Q_α) . Let us fix constants r_0, r_1 , $1 < r_0 < r_1$, a positive constant b , and a parameter $s \in (0, 1)$. By $\mathcal{H}(G_{(s)})$ we denote the space of functions which are holomorphic in $G_{(s)}$, where $G_{(s)} := \{z \in \mathbb{C} : |z| < r_0 + s(r_1 - r_0)\}$.

$$\mathbf{B} := \left\{ g = g(z, t) \in \bigcup_{0 < s < 1} \mathcal{C}([0, b(1-s)), \mathcal{H}(G_{(s)}) \cap \mathcal{C}^{1,\lambda}(\overline{G}_{(s)})) \right\} :$$

$$\|g\|_{\mathbf{B}} = \max \left\{ \sup_{s \in (0,1), h < b(1-s)} \max_{t \in [0,h]} \|g(\cdot, t)\|_{\mathcal{C}^\lambda(\overline{G}_{(s)})}; \right. \\ \left. \sup_{s \in (0,1), t < b(1-s)} \left\| \frac{\partial g}{\partial z}(\cdot, t) \right\|_{\mathcal{C}^\lambda(\overline{G}_{(s)})} \left(1 - \frac{t}{b(1-s)} \right)^{\frac{1}{2}} \right\} < \infty \}.$$

The problem (18) can be rewritten as:

$$\begin{aligned} \mathbf{I}(\phi) = \mathbf{I}(\phi_0) &+ \mathbf{A} \circ \mathbf{M}\left(\phi, \mathbf{M}(z, \mathbf{T}_t(\phi, \omega))\right) \\ &- 2\mathbf{J} \circ \mathbf{M}\left(\phi, \partial_z \circ \mathbf{M}(z, \mathbf{T}_t(\phi, \omega))\right), \end{aligned} \quad (21)$$

$$\mathbf{M}(\phi, \omega) := \phi \cdot \omega, \quad \mathbf{J}\psi := \int_0^t \psi(\cdot, \tau) d\tau.$$

$$\mathbf{B}_a := \left\{ g = g(z, t) \in \bigcup_{0 < s < 1} \mathcal{C}\left([0, b(1-s)), \mathcal{H}(A_{(s)}) \cap \mathcal{C}^{1,\lambda}(\overline{A}_{(s)})\right) : \right.$$

$$\|g\|_{\mathbf{B}_a} = \max \left\{ \sup_{s \in (0,1), h < b(1-s)} \max_{t \in [0,h]} \|g(\cdot, t)\|_{\mathcal{C}^\lambda(\overline{A}_{(s)})}; \right.$$

$$\left. \sup_{s \in (0,1), t < b(1-s)} \left\| \frac{\partial g}{\partial z}(\cdot, t) \right\|_{\mathcal{C}^\lambda(\overline{A}_{(s)})} \left(1 - \frac{t}{b(1-s)}\right)^{\frac{1}{2}} \right\} < \infty \},$$

$$A_{(s)} = \left\{ z \in \mathbb{C} : \frac{1}{r_0 + (r_1 - r_0)s} < |z| < r_0 + (r_1 - r_0)s \right\}, \quad 0 < s < 1.$$

Existence result. For each $\omega \in \mathbf{B}$ there exists a constant $b = b(\omega)$ such that the Cauchy-Kovalevsky problem (18) has a unique solution $\phi \in \mathbf{B}$, holomorphic and non-vanishing function in G_d . We prove that ω connected with ϕ by

$$\operatorname{Im}(\omega(z)) = \alpha \partial_{\theta} \left(|\phi(z)| \left(Q + \operatorname{Re}(\omega(z)) \right) \right) \text{ on } \partial G_1, \quad (22)$$

$$\omega(0) = 0, \quad (23)$$

is also holomorphic in G_d . Let $\phi \in \mathcal{H}(G_1) \cap \mathcal{C}^{1,\lambda}(\overline{G}_1)$, $\phi(z) \neq 0$. Then problem (22)–(23) possesses a unique solution $\omega(z) \in \mathcal{H}(G_1) \cap \mathcal{C}^{1,\lambda}(\overline{G}_1)$.

Main Theorem. *There exists an interval of time $[0, b)$ such that the problem (Q_α) has a unique solution (ϕ, ω) . $\phi = \phi(z, t)$ has no zeros on $\overline{G}_{r_0} \times [0, b)$ and belongs to the space $\mathcal{C}^1([0, b), \mathcal{H}(G_{r_0}) \cap \mathcal{C}^{1,\lambda}(\overline{G}_{r_0}))$. $\omega = \omega(z, t)$ belongs to the space $\mathcal{C}([0, b), \mathcal{H}(G_{r_0}) \cap \mathcal{C}^{1,\lambda}(\overline{G}_{r_0}))$. The constant r_0 is taken as in the definition of \mathbf{B} .*

Case 2. Hele-Shaw flow in a bounded domain. The polymer is injected through the gate AD under a constant pressure p_0 and occupied a domain Ω_t at the time t . The flow reaches the back wall along of arcs $\{(\alpha_k, \beta_k), k = 1, 2, \dots, n\}$. The arcs $\{(\gamma_k, \delta_k), k = 1, 2, \dots, n + 1\}$, where $\gamma_1 \in (A, B)$, $\gamma_k = \beta_{k-1}, k = 2, \dots, n + 1$, $\delta_k = \alpha_k, k = 1, \dots, n$, $\delta_{n+1} \in (C, D)$ is a free part $\Gamma_{f,t}$ of $\partial\Omega_t$, $\Gamma_{w,t}$ is the part of $\partial\Omega_t$ which reached by the polymer (i.e., $(A, \gamma_1) \bigcup_{k=1}^n (\alpha_k, \beta_k) \cup (\delta_{n+1}, D)$), and by Γ_* the gate (A, D) .

The flow satisfying Darcy law is described in terms $\omega = \varphi + i\psi$, $\varphi(z, t) \equiv \varphi(x, y, t)$ is a velocity potential, $\psi(z, t) \equiv \psi(x, y, t)$ is a stream function. It is assumed that the (impermeable) walls (A, B) , (D, C) , as well as the arcs (α_k, β_k) are stream lines, but the free arcs are equipotential lines. More precisely, $\psi|_{(A, B)} = \nu_1$, $\psi|_{(D, C)} = \nu_2$, $\psi|_{(\alpha_k, \beta_k)} = \nu, k = 1, 2, \dots, n$.

It leads

$$\operatorname{Re} \left[i \frac{\partial \omega}{\partial t} \right] = 0, \text{ on the wall.} \quad (24)$$

On the gate $(\overset{\frown}{A}, D)$ we have a constant pressure $p = p_0$, wlog:

$$\varphi = 0, \text{ on the free boundary.} \quad (25)$$

Besides,

$$\frac{\partial \varphi}{\partial n} = \frac{\partial n}{\partial t}, \text{ on the free boundary.} \quad (26)$$

Combining (25) and (26) one can get

$$\frac{\partial \varphi}{\partial t} = - \left(\frac{\partial n}{\partial t} \right)^2, \text{ on the free boundary,}$$

or in terms of the complex potential:

$$\operatorname{Re} \omega = 0, \operatorname{Re} \left[\frac{\partial \omega}{\partial t} \right] = - \left| \frac{\partial \omega}{\partial z} \right|^2, \text{ on the free boundary.} \quad (27)$$

Hele-Shaw flow in a bounded cell. Mathematical model. Let $f(\zeta, t)$ be a conformal mapping of the unit disc G_1 onto domain Ω_t , $W(\zeta, t)$ be analytic in the unit disc G_1 , s.t.:

$$\operatorname{Re} \left\{ \overline{\Lambda(\tau)} W(\tau, t) \right\} = c(\tau) = \begin{cases} 0, & \tau \in \tilde{\Gamma}_{f,t}, \\ \nu, & \tau \in \bigcup_{k=1}^n (\tilde{\alpha}_k, \tilde{\beta}_k), \\ \nu_1, & \tau \in (\tilde{A}, \tilde{\gamma}_1), \\ \nu_2, & \tau \in (\tilde{\delta}_{n+1}, \tilde{D}), \\ p_0, & \tau \in (\tilde{A}, \tilde{D}), \end{cases} \quad (28)$$

$$\Lambda(\tau) = \begin{cases} 1, & \tau \in \tilde{\Gamma}_{f,t}, \\ -i, & \tau \in \tilde{\Gamma}_{w,t}, \\ 1, & \tau \in \tilde{\Gamma}_*. \end{cases} \quad (29)$$

$$\omega(z, t) = W \left[f^{-1}(z, t), t \right]. \quad (30)$$

Then the problem is equivalent to the Riemann-Hilbert problem

$$\operatorname{Re} \left\{ \overline{\Lambda(\tau)} \Phi(\tau, t) \right\} = d(\tau, t) = \begin{cases} 0, & \tau \in \tilde{\Gamma}_{w,t} \cup \tilde{\Gamma}_*, \\ \left| \frac{W'(\tau, t)}{f'(\tau, t)} \right|^2, & \tau \in \tilde{\Gamma}_{f,t}, \end{cases} \quad (31)$$

$$\Phi(\zeta, t) := W'(\zeta, t) \frac{\dot{f}(\zeta, t)}{f'(\zeta, t)}. \quad (32)$$

$$\operatorname{Re} \left\{ \overline{\Lambda(\tau)} W(\tau) \right\} = c(\tau), \quad \tau \in \mathbb{T} \setminus \{\tau_1^t, \dots, \tau_{2n+4}^t\}, \quad (33)$$

$$\operatorname{Re} \left\{ \overline{\Lambda(\tau)} W'(\tau) \frac{\dot{f}(\tau, t)}{f'(\tau, t)} \right\} = d(\tau, t), \quad \tau \in \partial G_1 \setminus \{\tau_1^t, \dots, \tau_{2n+4}^t\}, \quad (34)$$

$$f(\zeta, 0) = f_0(\zeta), \quad \zeta \in \mathbb{U}. \quad (35)$$

The problem (33) has the form

$$\operatorname{Re} \left\{ \overline{\Lambda(\tau)} g(\tau) \right\} = h(\tau), \quad \tau \in \partial G_1 \setminus \{\tau_1^t, \dots, \tau_{2n+4}^t\}, \quad (36)$$

$$\varkappa = n + 2;$$

$$W(\zeta) = W_p(\zeta) + X(\zeta)P_\varkappa(\zeta),$$

$$X(\zeta) = C \left(\prod_{j=1}^{2n+4} (\zeta - \tau_j^t) \right)^{-1/2},$$

$$P_\varkappa(\zeta) = c_0 \zeta^\varkappa + c_1 \zeta^{\varkappa-1} + \dots + c_\varkappa, \quad c_k = \overline{c_{\varkappa-k}}, \quad k = 0, 1, \dots, \varkappa,$$

$$W_p(\zeta) = \frac{X(\zeta)}{2\pi i} \int_{\mathbb{T}} \frac{\tau^{\varkappa+1} + \zeta^{\varkappa+1} h(\tau) d\tau}{X^+(\tau) \tau^{\varkappa+1} (\tau - \zeta)}.$$

By using the solution operator \mathbf{T}_Λ^t to (36) we rewrite the boundary condition (34) in the following equivalent form

$$f \frac{W'}{f'} = \mathbf{T}_\Lambda^t(d), \quad \zeta \in \overline{G_1} \setminus \{\tau_1^t, \dots, \tau_{2n+4}^t\}. \quad (37)$$

By using new unknown function

$$\phi(\zeta, t) := \frac{W'(\zeta, t)}{f'(\zeta, t)}, \quad \zeta \in \overline{G_1} \setminus \{\tau_1, \dots, \tau_{2n+4}\}, \quad (38)$$

we rewrite the relations (34), (35) in the form of an abstract Cauchy-Kovalevsky problem

$$\left\{ \begin{array}{l} \frac{\partial \phi}{\partial t}(\zeta, t) = \frac{1}{W'(\zeta, t)} \frac{\partial \phi}{\partial \zeta}(\zeta, t) \mathbf{T}_\Lambda^t(d)(\zeta, t) \\ \quad - \frac{1}{W'(\zeta, t)}(\zeta) \phi(\zeta, t) \frac{\partial}{\partial \zeta} \mathbf{T}_\Lambda^t(d)(\zeta, t), \quad (\zeta, t) \in \partial G_1 \times [0, T], \\ \phi(\zeta, 0) = \frac{W'(\zeta, 0)}{(f_0)'(\zeta)}, \quad \zeta \in G_1. \end{array} \right. \quad (39)$$

Let us fix numbers $\alpha_k \in \mathbb{C}$ and define the weight functions

$$\rho^{(j)}(t) = \prod_{k=1}^m |t - t_k|^{\alpha_k + j}, \quad j \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}, \quad (40)$$

$$\varphi \in \mathbb{A}_{\mu;n}^r(\Gamma; t_1, \dots, t_m) \Leftrightarrow \begin{cases} \varphi \Big|_{\overset{\smile}{\tau_k, \tau_{k+1}}} \in \mathbb{B}_{\mu}^r(t_k, t_{k+1}), & k = 1, 3, \dots, \\ \varphi \Big|_{\overset{\smile}{\tau_k, \tau_{k+1}}} \in \mathbb{KZ}_{n+\mu}^0(t_k, t_{k+1}), & k = 2, 4, \dots \end{cases} \quad (41)$$

The conditions on the arcs with even indexes ($k = 2, 4, \dots$):

$$\begin{aligned} \varphi \in \mathbb{KZ}_{n+\mu}^0(t_k, t_{k+1}) &\Leftrightarrow \varphi_j := \rho^{(j)} \partial^j \varphi \in \mathbb{Z}_{\mu}^0(t_k, t_{k+1}) \quad (j = 0, 1, \dots, n) \\ &\Leftrightarrow \varphi_j \in \mathbb{Z}_{\mu}(t_k, t_{k+1}); \varphi_j(t_k) = \varphi_j(t_{k+1}) = 0. \end{aligned} \quad (42)$$

For each α_k the expression $\rho^{(j)} \partial^j \varphi$ is defined by the equality

$$\rho^{(j)} \partial^j \varphi := \rho^{(j)}(\omega(\tau)) \frac{\partial^j \varphi(\omega(\tau))}{\partial \tau^j}. \quad (43)$$

The spaces $\mathbb{Z}_\mu(t_k, t_{k+1})$, $\mathbb{Z}_\mu^0(t_k, t_{k+1})$:

$$\|\varphi | \mathbb{Z}_\mu^0(t_k, t_{k+1})\| = \|\rho\varphi | \mathbb{Z}_\mu(t_k, t_{k+1})\|, \quad (44)$$

$$\|\psi | \mathbb{Z}_\mu(t_k, t_{k+1})\| := \sup_{\tau \in (\tau_k, \tau_{k+1})} |\psi(\omega(\tau))| + \quad (45)$$

$$+ \sup_{\tau \in (\tau_k, \tau_{k+1}); h > 0} h^{-\mu} |\psi(\omega(\tau + h)) - 2\psi(\omega(\tau)) + \psi(\omega(\tau - h))|.$$

The norm in the space $\mathbb{KZ}_{n+\mu}^0(t_k, t_{k+1})$ is defined by the standard way:

$$\left\| \varphi \middle| \mathbb{KZ}_{n+\mu}^0(t_k, t_{k+1}) \right\| := \sum_{j=0}^{n-1} \sup_{\tau \in (\tau_k, \tau_{k+1})} |\varphi_j(\omega(\tau))| + \left\| \varphi_n \middle| \mathbb{Z}_{\mu}^0(t_k, t_{k+1}) \right\|. \quad (46)$$

$$\partial^j \varphi = o\left(\rho^{(j)}\right), \quad t \rightarrow t_k \quad \text{or} \quad t \rightarrow t_{k+1}. \quad (47)$$

It will be said that $\varphi \in \mathbb{B}_{\mu}^r(t_k, t_{k+1})$, $k = 1, 3, \dots$, if $\exists r > 0$ such that

$$\varphi \in \mathcal{C}^{\infty}(t_k, t_{k+1}) : \sup_{j \in \mathbb{N}_0} \left\| \varphi_j \middle| \mathbb{Z}_{\mu}(t_k, t_{k+1}) \right\| \cdot \frac{r^j}{\Gamma_1(j)} < +\infty. \quad (48)$$

The last supremum determines the norm in the space $\mathbb{B}_{\mu}^r(t_k, t_{k+1})$:

$$\left\| \varphi_j \middle| \mathbb{B}_{\mu}^r(t_k, t_{k+1}) \right\| := \sup_{j \in \mathbb{N}_0} \left\| \rho^{(j)} \partial^{(j)} \varphi \middle| \mathbb{Z}_{\mu}(t_k, t_{k+1}) \right\| \cdot \frac{r^j}{\Gamma_1(j)}. \quad (49)$$

Scale of Banach spaces. We choose $\mathbb{A}_{\mu;n}^r(\text{cl } D; t_1, \dots, t_m)$ as a subspace of n -times piece-wise continuously differentiable functions.

$$\left\| \varphi \Big|_{\mathbb{A}_{\mu;n}^r(\Gamma; t_1, \dots, t_m)} \right\| := \max \left\{ \max_{k=1,3,\dots} \left\| \varphi \Big|_{\overset{\smile}{t_k, t_{k+1}}} \Big|_{\mathbb{B}_{\mu}^r(t_k, t_{k+1})} \right\|, \max_{k=2,4,\dots} \left\| \varphi \Big|_{\overset{\smile}{t_k, t_{k+1}}} \Big|_{\mathbb{KZ}_{n+\mu}^0(t_k, t_{k+1})} \right\| \right\}. \quad (50)$$

$$\mathbb{A}_{\mu;n}^r(\text{cl } D; t_1, \dots, t_m) := \mathcal{H}(D) \cap \mathbb{A}_{\mu;n}^r(\Gamma; t_1, \dots, t_m). \quad (51)$$

$$\mathbb{A}_s(\Gamma_{(s)}) := \mathbb{A}_{\mu;n}^r(\Gamma_{(s)}; t_1(s), \dots, t_m(s)). \quad (52)$$

Introduce also the space of functions defined in the domains:

$$\mathbb{A}_s(\text{cl } D_{(s)}) := \mathcal{H}(D_{(s)}) \cap \mathbb{A}_s(\Gamma_{(s)}). \quad (53)$$

Let us fix a positive number r and numbers $1 < r_0 < r_1$.

$$D_{\{s\}} := D_{(r_0+s(r_1-r_0))}, \Gamma_{\{s\}} := \partial D_{\{s\}}. \quad (54)$$

$t_1^{(s)}, t_2^{(s)}, \dots, t_m^{(s)}$ - knots. $\mathbb{A}_{\{s\}}(\text{cl } D_{\{s\}}) := \mathbb{A}_{r_0+s(r_1-r_0)}(\text{cl } D_{(r_0+s(r_1-r_0))})$,
 $\mathbb{A}_{\{s\}}(\Gamma_{\{s\}}) := \mathbb{A}_{r_0+s(r_1-r_0)}(\Gamma_{(r_0+s(r_1-r_0))})$. The scale of Banach spaces

$$\left\| \varphi \Big|_{\mathbb{B}_{s,p}(D_{\{s\}}; \Gamma_1; \rho)} \right\| := \max \left\{ \sup_{z \in D_{\{s\}}} |\varphi(z)|, \max_{k=1,3,\dots} \left\| \varphi \Big|_{\overset{\smile}{t_k^{(s)}, t_{k+1}^{(s)}}} \Big|_{\mathbb{B}_\mu^{pr}(t_k^{(s)}, t_{k+1}^{(s)})} \right\|, \right.$$

$$\left. \max_{k=2,4,\dots} \left\| \varphi \Big|_{\overset{\smile}{t_k^{(s)}, t_{k+1}^{(s)}}} \Big|_{\mathbb{KZ}_{n+\mu}^0(t_k^{(s)}, t_{k+1}^{(s)})} \right\| \right\} < +\infty.$$

(55)

$$\mathbf{B}_{pw,n} := \left\{ g(\zeta, t) \in \bigcup_{0 < s < 1, 0 < p < 1} \mathcal{C}([0, \delta(1-s)), \mathbb{B}_{s,p}(D_{\{s\}}; \Gamma_1; \rho)) : (56) \right.$$

$$\|g\|_{\mathbf{B}_{pw,n}} := \max \left\{ \sup_{s \in (0,1), d < \delta(1-s)} \max_{t \in [0,d]} \sup_{z \in \mathbb{D}_{\{s\}}} \|g(z, t)\|, \right.$$

$$\sup_{s \in (0,1), p \in (0,1), d < \delta(1-s)} \max_{k=1,3,\dots,2n+3} \max_{t \in [0,d]} \left\| g(\cdot, t) \Big|_{\tau_k^{(s)}, \tau_{k+1}^{(s)}} \Big|_{\mathbb{B}_{\mu}^{pr}(\tau_k^{(s)}, \tau_{k+1}^{(s)})} \right\|,$$

$$\sup_{s \in (0,1), d < \delta(1-s)} \max_{k=2,4,\dots,2n+2} \max_{t \in [0,d]} \left\| g(\cdot, t) \Big|_{\tau_k^{(s)}, \tau_{k+1}^{(s)}} \Big|_{\mathbb{KZ}_{1+\mu}^0(\tau_k^{(s)}, \tau_{k+1}^{(s)})} \right\|,$$

$$\sup_{s \in (0,1), p \in (0,1), d < \delta(1-s)} \max_{t \in [0,d]} \left\| \rho^{(1)}(\cdot) \frac{\partial g}{\partial \zeta}(\cdot, t) \Big|_{\mathbb{Z}_{\{\mu\}}(\overline{D}_{\{s\}})} \left(1 - \frac{t}{\delta(1-s)} \right)^{1/2} \right\| < \infty \Big\}$$

Main Theorem. Let $\nu, \nu_1, \nu_2, p_0 > 0$ be positive constants. Let $f_0 = f_0(\zeta)$ maps G_1 onto Ω_0 , and can be continued for the function analytic and univalent in a neighborhood $G_1 \setminus \{\tau_1^0, \dots, \tau_{2n+4}^0\} \equiv G_1 \bigcup_{k=1}^{n+1} U(\tau_{2k}^0, \tau_{2k+1}^0)$,

$$U(\tau_{2k}^0, \tau_{2k+1}^0) \cap \text{cl}[\tau_{2k}^0, \tau_{2k+1}^0] = (\tau_{2k}^0, \tau_{2k+1}^0).$$

$$t_k^0 \equiv f_0(\tau_k^0), \quad k = 1, \dots, 2n + 4,$$

be the images of τ_k^0 under the mapping f_0 .

Then for each $t \in [0, \delta)$ with δ sufficiently small there exists a unique solution $\omega(z, t) \equiv W[f^{-1}(z, t), t]$ to (33), (34), (35) in the space

$$C^1 \left([0, \delta), \mathcal{H}(D_{\{0\}}) \cap \left(\bigcup_{k=1}^{n+1} \mathbb{KZ}_{2+\mu}^0(\tau_{2k}^{(0)}, \tau_{2k+1}^{(0)}) \right) \cap \left(\bigcup_{k=1}^{n+2} \mathbb{B}_\mu^{pr}(\tau_{2k-1}^{(0)}, \tau_{2k}^{(0)}) \right) \right).$$

Real-variable Hele-Shaw model.

Viscous incompressible fluid occupies a doubly connected domain $D_1(t)$ at a time instant $t \geq 0$. Internal domain F is a fixed small obstacle (hole). The simply connected domain without hole will be denoted $D(t)$, $\delta := \text{diam } F > 0$, $D(0)$ is open bounded set with a smooth boundary

$$\text{dist } \{\partial F, \partial D(0)\} = 2d > 0. \quad (57)$$

$$\partial F \in \mathcal{C}^{2,\alpha}, \quad \partial D(0) \in \mathcal{C}^{1,\alpha}, \quad 0 < \alpha < 1. \quad (58)$$

2D potential flow of incompressible fluid in the Hele-Shaw cell:

$$\frac{\partial \mathbf{V}}{\partial t} = 0, \quad V_3 = 0, \quad (59)$$

$$\mathbf{V} = -\frac{h^2}{12\mu} \nabla p. \quad (60)$$

$$\Delta p = 0. \quad (61)$$

Unique source/sink at a fixed point $z_0 = (x_0, y_0)$

$$z_0 \in D_1(0) = D(0) \setminus cl F, \quad O \in int F.$$

$$p(z, t) \sim -\frac{Q(t)}{2\pi} \log |z - z_0|, \quad |z| \rightarrow z_0. \quad (62)$$

$$\frac{\partial p}{\partial n} = 0, \quad z \in \partial F. \quad (63)$$

$$p(z) = 0, \quad z \in \Gamma(t), \quad (64)$$

$$\frac{d\Gamma}{dt} = \mathbf{V}, \quad z \in \Gamma(t). \quad (65)$$

$$\frac{d\Gamma}{dt} = -\frac{h^2}{12\mu} \nabla p. \quad (66)$$

Let us now introduce new unknown function, a one-parametric family of \mathcal{C}^2 -diffeomorphisms

$$w(s, t) = (u(s, t), v(s, t)) : \partial G_1 \times I \rightarrow \Gamma(t), \quad G_1 = \{s = (s_1, s_2) \in \mathbb{R}^2 : |s| < 1\}. \quad (67)$$

The function $w(s, t)$ in (67) determines an unknown parametrization of the free boundary $\Gamma(t)$, namely,

- (i) $w(s, t) \in \Gamma(t)$ for all $(s, t) \in \partial G_1 \times I$,
- (ii) $w(\cdot, t) : \partial G_1 \rightarrow \Gamma(t)$ is a \mathcal{C}^2 -diffeomorphism for each fixed $t \in I$,
- (iii) $w(\cdot, \cdot) \in \mathcal{C}^2(\partial G_1 \times I; \mathbb{R}^2)$.

$$p = Q \cdot \mathcal{G}_{D_1}(t), \quad (68)$$

and $\mathcal{G}_{D_1(t)}$ is the solution to

$$\Delta \mathcal{G}_{D_1(t)}(z, z_0) + \delta_0(z - z_0) = 0, \quad z \in D_1(t), \quad (69)$$

$$\mathcal{G}_{D_1(t)}(z, z_0) = 0, \quad z \in \Gamma(t), \quad (70)$$

$$\frac{\partial \mathcal{G}_{D_1(t)}}{\partial n}(z, z_0) = 0, \quad z \in \partial F. \quad (71)$$

Problem (HS₀). Find a pair $\{w(s,t); \mathcal{G}(z, z_0; t)\}$, such that $w(s,t) : \partial G_1 \times I \rightarrow \mathbb{R}^2$ is a \mathcal{C}^2 -diffeomorphism satisfying

(i) $w(s,t) \in \Gamma(t)$ for all $(s,t) \in \partial G_1 \times I$;

(ii) $w(\cdot, t) : \partial G_1 \rightarrow \Gamma(t)$ is a \mathcal{C}^2 -diffeomorphism for each fixed $t \in I$;

(iii) $w^{(0)}(s) = w(s, 0)$ is a given \mathcal{C}^2 -diffeomorphism of the unit circle ∂G_1 , which describes the boundary $\Gamma(0)$ of initial domain $D_1(0)$;

(iv) $\mathcal{G}(z, z_0; t)$ is Green's function of the operator $-\Delta$ in the doubly connected domain $D_1(t)$ with the homogeneous Neumann data on the fixed boundary ∂F and the homogeneous Dirichlet data on the free boundary $\Gamma(t)$, i.e. satisfies conditions (69)–(71) for each fixed $t \in I$;

(v) $\partial_t w(s,t) = -\frac{Qh^2}{12\mu} \cdot \nabla \mathcal{G}(w(s,t), z_0; t)$ for all $(s,t) \in \partial G_1 \times I$.

We consider four essentially different situations:

(a) the source z_0 and ALL points $z = w(s, t)$ of the boundary $\partial\Gamma(0)$ of the initial domain $D_1(0)$ are distant from the boundary $\partial\omega_\varepsilon$ of the obstacle;

(b) the source z_0 is close to the boundary $\partial\omega_\varepsilon$ of the obstacle, but ALL points $z = w(s, t)$ of $\partial\Gamma(0)$ are distant from $\partial\omega_\varepsilon$;

(c) the source z_0 and SOME points $z = w(s, t)$ of $\partial\Gamma(0)$ are close to the boundary $\partial\omega_\varepsilon$ of the obstacle; we consider those points $z = w(s, t)$ which are distant from $\partial\omega_\varepsilon$;

(d) the source z_0 and SOME points $z = w(s, t)$ of $\partial\Gamma(0)$ are close to the boundary $\partial\omega_\varepsilon$ of the obstacle; we consider those points $z = w(s, t)$ which are close to $\partial\omega_\varepsilon$.

1. B. Gustafsson, A. Vasil'ev, Conformal and Potential Analysis in Hele-Shaw cells, Birkhäuser Verlag, Basel-Boston-Berlin, 2006.
2. V. Maz'ya, A. Movchan, Uniform asymptotics of Green's kernels for mixed and Neumann problems in domains with small holes and inclusions. Isakov, Victor (ed.), Sobolev spaces in mathematics. III: Applications in mathematical physics. New York, NY: Springer; Novosibirsk: Tamara Rozhkovskaya Publisher. International Mathematical Series, 10, 2009, 277–316.
3. V. Maz'ya, A. Movchan, Uniform asymptotics of Green's kernels in perforated domains and meso-scale approximation // Complex Var. Elliptic Equ., 57, No. 2-4, 137-154 (2012).
4. V. Maz'ya, A. Movchan, M. Nieves, Green's kernels and meso-scale approximations in perforated domains. Lecture Notes in Mathematics 2077. Berlin: Springer. xvii, 254 p. (2013).
5. M. Reissig, S. Rogosin, with an appendix of F. Hübner, Analytical and numerical treatment of a complex model for Hele-Shaw moving boundary value problems with kinetic undercooling regularization, *Euro J. Appl. Math.*, **10** (1999), 561–579.
6. S.V. Rogosin, Hele-Shaw moving boundary value problems in a bounded domain. Local in time solvability, *Complex Variables*, **50**, No. 7-11 (2005), 745–764.