Coalescing random walkers

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Avalanche size in the totally asymmetric Oslo model maps to the area inscribed by annihilating random walkers.



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Scaling arguments



Fluctuations *f*, time *t*, diffusion constant *D* (Brownian motion!), initial gap x_0 . Characteristic area $s_c \propto ft \propto (Dt)^{1/2} t \propto t^{3/2}$



First moment: N walkers cover area Nx_0t plus error independent of N, when large enough.

Surprise: Error vanishes like 1/N, rather than $1/\sqrt{N}$ for independent areas.

This latter term is cancelled by short-lived anticorrelations, as seen in the variance of the estimated average area:

$$\frac{1}{N^2}\sum_{ij}^{N}\left(\left\langle s_is_j\right\rangle - \left\langle s_i\right\rangle\left\langle s_j\right\rangle\right) = \frac{\sigma^2(s)}{N} + \sum_{i,j=1\atop l\neq j}^{N}\left(\left\langle s_is_j\right\rangle - \left\langle s_i\right\rangle\left\langle s_j\right\rangle\right) = \frac{2}{3N^2}Dt^3 + \dots$$





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Scaling arguments



- Assume simple scaling of cluster size distribution $P(s) = s^{-\tau} \mathcal{G}(s/t^{d_f}).$
- Moments thus go like $\langle s^n \rangle \propto t^{d_f(1-\tau+n)}$.
- Characteristic size scales like $s_c \propto \sqrt{D} t^{3/2}$, *i.e.* $d_f = 3/2$.
- First moment goes like $\langle s \rangle = x_0 t$, so that $\tau = 4/3$.
- Higher moments thus go like $\langle s^n \rangle = A_n x_0 D^{(1-\tau+n)/2} t^{(3n-1)/2}$.
- All moments (precisely: their leading orders in t) linear in x₀?!?



$$\langle \boldsymbol{s}^n \rangle = (\boldsymbol{x}_0 t)^n \qquad \langle \boldsymbol{s}^n \rangle = \boldsymbol{A}_n \boldsymbol{x}_0 \boldsymbol{D}^{(1-\tau+n)/2} t^{(3n-1)/2}$$



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As long as characteristic size does not involve x_0 , all moments are linear in x_0 .

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Perturbation theory

Calculation of moments in closed form is difficult.

Perturbation theory in powers of t actually not possible (integrals violate assumptions about small and large parameters). Overlooking that and ignoring divergent coefficients, gives

$$\langle s^2 \rangle = c_0 t^{5/2} + c_1 t^{3/2} + c_2 t^{1/2} + c_3 t^{-1/2} + \dots$$

Series expansion of the exact result:

$$\langle s^2 \rangle = c_0 t^{5/2} + c_1 t^{3/2} - \mathbf{d_0} \mathbf{t} + c_2 t^{1/2} - \mathbf{d_1} + c_3 t^{-1/2} + \dots$$

Exact results

Exact second moment is messy:

$$\begin{split} \left\langle s^{2} \right\rangle &= \frac{1}{180} \left(\frac{x_{0}^{3}}{D} \right)^{2} \left[\left(1 + 28 \frac{tD}{x_{0}^{2}} + 132 \left(\frac{tD}{x_{0}^{2}} \right)^{2} \right) \sqrt{\frac{4tD}{\pi x_{0}^{2}}} \exp\left(-\frac{x_{0}^{2}}{4tD} \right) \right. \\ &\left. - \left(1 + 30 \frac{tD}{x_{0}^{2}} \right) \operatorname{erfc}\left(\sqrt{\frac{x_{0}^{2}}{4tD}} \right) \right. \\ &\left. + 60 \left(\frac{tD}{x_{0}^{2}} \right)^{2} \left(3 + 2 \frac{tD}{x_{0}^{2}} \right) \operatorname{erf}\left(\sqrt{\frac{x_{0}^{2}}{4tD}} \right) \right] \end{split}$$

Correlation functions

How to calculate the correlation function $\langle s_i s_{i+1} \rangle$?

- (Idea by Alan Bray)
 - Observation:

$$\langle (s_i + s_{i+1})^2 \rangle = \langle s_i^2 \rangle + \langle s_{i+1}^2 \rangle + 2 \langle s_i s_{i+1} \rangle.$$

- $\langle s_i^2 \rangle = \langle s_{i+1}^2 \rangle$ known from above.
- $\langle (s_i + s_{i+1})^2 \rangle$ is the second moment of the area for the initial gap *doubled*, $x_0 \rightarrow 2x_0$.
- Thus:

$$\langle s_i s_{i+1} \rangle = \frac{1}{2} \left\langle s^2 \right\rangle (2x_0) - \left\langle s^2 \right\rangle (x_0) \;.$$

• Similar for higher correlation functions $\langle s_i s_{i+j} \rangle$. Note: Leading order terms, linear in x_0 , cancel.

Summary

Area between coalescing random walkers.

- Scaling arguments suggest the right behaviour.
- Anti-correlations visible when estimating first moment.
- All moments linear in initial gap (to leading order in *t*).
- Perturbation theory fails.
- First and second moment known exactly.
- Correlation function can be calculated by merging areas.

Published in:

Peter Welinder, Gunnar Pruessner and Kim Christensen *Multiscaling in the sequence of areas enclosed by coalescing random walkers*, New J. Phys. **9**, 149-1–18 (2007).

Thank you!